

Observational consequences of charged Planck mass gravitini

Krzysztof A. Meissner

**FACULTY OF
PHYSICS**
UNIVERSITY
OF WARSAW

AEI, HermannFest, 14.09.2022



GoIm 2022

Common papers

1. K.A. Meissner and H. Nicolai, *Conformal Symmetry and the Standard Model*, Phys.Lett. **B648** (2007) 312
2. K.A. Meissner and H. Nicolai, Phys. Lett. **B660** (2008) 260
3. K.A. Meissner and H. Nicolai, Eur.Phys.J. **C57** (2008) 493
4. K.A. Meissner and H. Nicolai, Acta Phys.Pol. **B40** (2009) 2737
5. K.A. Meissner and H. Nicolai, Phys.Rev. **D80:086005** (2009)
6. A. Latosiński, K.A. Meissner and H. Nicolai, arXiv:1010.5417 [hep-ph]
7. A. Latosiński, K.A. Meissner and H. Nicolai, Nucl. Phys. **B868** (2013) 596-626
8. A. Latosiński, K.A. Meissner and H. Nicolai, Eur.Phys.J. **C73** (2013) 2336
9. K.A. Meissner and H. Nicolai, Phys. Lett **B718** (2013), 943
10. P. Dutta, K.A. Meissner and H. Nicolai, Phys. Rev. **D87** (2013) 105019
11. P.H. Chankowski, A. Lewandowski, K.A. Meissner and H. Nicolai, Mod. Phys. Lett. **A 30** No. 2 (2015) 1550006
12. K.A. Meissner and H. Nicolai, Phys. Rev. **D91** (2015) 065029
13. A.Latosiński, A. Lewandowski, K.A. Meissner and H. Nicolai, JHEP **1510** (2015) 170

Common papers

14. K.A. Meissner and H. Nicolai, Phys. Lett. **B772** (2017) 169
15. H. Godazgar, K.A. Meissner and H. Nicolai, JHEP **04** (2017) 165
16. K.A. Meissner and H. Nicolai, Phys.Rev. **D96** (2017) 041701
17. A. Lewandowski, K.A. Meissner and H. Nicolai, Phys. Rev. **D97** (2018) 035024.
18. K.A. Meissner and H. Nicolai, Foundations of Physics **48** (2018) 1150
19. G.F.R. Ellis, K.A. Meissner and H. Nicolai, Nature Physics **14** (2018) 770
20. K.A. Meissner, H. Nicolai and J. Plefka, Phys. Lett. **B791** (2019) 62
21. K.A. Meissner and H. Nicolai, *Standard Model Fermions and Infinite-Dimensional R-Symmetries*, Phys. Rev. Lett. **121** (2018) 091601
22. K.A. Meissner and H. Nicolai, *Planck Mass Charged Gravitino Dark Matter*, Phys.Rev. **D100** (2019) 035001
23. K.A. Meissner and H. Nicolai, *Superheavy Gravitinos and Ultra-High Energy Cosmic Rays* JCAP **09** (2019) 041
24. K.A. Meissner and H. Nicolai, *Supermassive gravitinos and giant primordial black holes*, Phys.Rev. **D102** (2020) 103008
25. K.A. Meissner and H. Nicolai, *Origin and growth of primordial black holes*, Phys.Lett. **B819** (2021) 136468

$N = 8$ supergravity

E. Cremmer, B. Julia, J.Scherk, B. de Wit, H. Nicolai,...

- it is very special – maximal, has hidden symmetries...

$N = 8$ supergravity

E. Cremmer, B. Julia, J.Scherk, B. de Wit, H. Nicolai,...

- it is very special – maximal, has hidden symmetries...
- it is conformal anomaly free

K.A.M., H. Nicolai, Phys.Lett.B 772 (2017) 169

$N = 8$ supergravity

E. Cremmer, B. Julia, J.Scherk, B. de Wit, H. Nicolai,...

- it is very special – maximal, has hidden symmetries...
- it is conformal anomaly free

K.A.M., H. Nicolai, Phys.Lett.B 772 (2017) 169

- $N = 8$ supergravity field content:
 - 1 graviton
 - 8 gravitinos
 - 28 vectors
 - 56 spin 1/2 fermions
 - 70 scalars

$N = 8$ supergravity

E. Cremmer, B. Julia, J.Scherk, B. de Wit, H. Nicolai,...

- it is very special – maximal, has hidden symmetries...
- it is conformal anomaly free

K.A.M., H. Nicolai, Phys.Lett.B 772 (2017) 169

- $N = 8$ supergravity field content:
 - 1 graviton
 - 8 gravitinos
 - 28 vectors
 - 56 spin 1/2 fermions
 - 70 scalars
- we assume that there is some truth in it, at least in field content (maybe as some limit of M theory) even if we don't expect that it is realized at any scale

$N = 8$ SUGRA and the Standard Model

M. Gell-Mann, H. Nicolai, N. Warner

- after full susy breaking it has exactly 48 'massless' fermion dofs χ^{ijk} as in SM with 6 quarks and 6 leptons
– possible explanation of 3 generations!

$N = 8$ SUGRA and the Standard Model

M. Gell-Mann, H. Nicolai, N. Warner

- after full susy breaking it has exactly 48 'massless' fermion dofs χ^{ijk} as in SM with 6 quarks and 6 leptons
– possible explanation of 3 generations!
- the original Gell-Mann conjecture of $SU(3) \times U(1)$ (extended by Nicolai and Warner) gives proper group assignments for both quarks and leptons but the electric charges are shifted by $\pm 1/6$ ($SU(2)$ would play a very different role)

$N = 8$ SUGRA and the Standard Model

M. Gell-Mann, H. Nicolai, N. Warner

- after full susy breaking it has exactly 48 'massless' fermion dofs χ^{ijk} as in SM with 6 quarks and 6 leptons – possible explanation of 3 generations!
- the original Gell-Mann conjecture of $SU(3) \times U(1)$ (extended by Nicolai and Warner) gives proper group assignments for both quarks and leptons but the electric charges are shifted by $\pm 1/6$ ($SU(2)$ would play a very different role)
- **a correction** of the usual $U(1)$ generator in $SU(3) \times U(1)$

$$\mathcal{J} = (T \wedge \mathbf{1} \wedge \mathbf{1} + \mathbf{1} \wedge T \wedge \mathbf{1} + \mathbf{1} \wedge \mathbf{1} \wedge T + T \wedge T \wedge T), \quad \mathcal{J}^2 = -1$$

gives proper quantum numbers of all quarks and leptons!

K.A.M., H. Nicolai, Phys. Rev. D91 (2015) 065029

$N = 8$ SUGRA and the Standard Model

M. Gell-Mann, H. Nicolai, N. Warner

- after full susy breaking it has exactly 48 'massless' fermion dofs χ^{ijk} as in SM with 6 quarks and 6 leptons – possible explanation of 3 generations!
- the original Gell-Mann conjecture of $SU(3) \times U(1)$ (extended by Nicolai and Warner) gives proper group assignments for both quarks and leptons but the electric charges are shifted by $\pm 1/6$ ($SU(2)$ would play a very different role)
- **a correction** of the usual $U(1)$ generator in $SU(3) \times U(1)$

$$\mathcal{J} = (T \wedge \mathbf{1} \wedge \mathbf{1} + \mathbf{1} \wedge T \wedge \mathbf{1} + \mathbf{1} \wedge \mathbf{1} \wedge T + T \wedge T \wedge T), \quad \mathcal{J}^2 = -1$$

gives proper quantum numbers of all quarks and leptons!

K.A.M., H. Nicolai, Phys. Rev. D91 (2015) 065029

- this correction is outside of $N = 8$ SUGRA!

Massive Gravitini

K.A.M., H. Nicolai, Phys. Rev. Lett. 121 (2018) 091601

- The eight massive gravitini under $SU(3) \times U(1)_{em}$ as

$$\left(\mathbf{3}, \frac{1}{3}\right) \oplus \left(\bar{\mathbf{3}}, -\frac{1}{3}\right) \oplus \left(\mathbf{1}, \frac{2}{3}\right) \oplus \left(\mathbf{1}, -\frac{2}{3}\right)$$

where the charges include the \mathcal{J} correction

Massive Gravitini

K.A.M., H. Nicolai, Phys. Rev. Lett. 121 (2018) 091601

- The eight massive gravitini under $SU(3) \times U(1)_{em}$ as

$$\left(\mathbf{3}, \frac{1}{3}\right) \oplus \left(\bar{\mathbf{3}}, -\frac{1}{3}\right) \oplus \left(\mathbf{1}, \frac{2}{3}\right) \oplus \left(\mathbf{1}, -\frac{2}{3}\right)$$

where the charges include the \mathcal{J} correction

- all gravitini carry *fractional electric charges!*

Massive Gravitini

K.A.M., H. Nicolai, Phys. Rev. Lett. 121 (2018) 091601

- The eight massive gravitini under $SU(3) \times U(1)_{em}$ as

$$\left(\mathbf{3}, \frac{1}{3}\right) \oplus \left(\bar{\mathbf{3}}, -\frac{1}{3}\right) \oplus \left(\mathbf{1}, \frac{2}{3}\right) \oplus \left(\mathbf{1}, -\frac{2}{3}\right)$$

where the charges include the \mathcal{J} correction

- all gravitini carry *fractional electric charges!*
- complex triplet of gravitini is in addition strongly interacting

Massive Gravitini

K.A.M., H. Nicolai, Phys. Rev. Lett. 121 (2018) 091601

- The eight massive gravitini under $SU(3) \times U(1)_{em}$ as

$$\left(\mathbf{3}, \frac{1}{3}\right) \oplus \left(\bar{\mathbf{3}}, -\frac{1}{3}\right) \oplus \left(\mathbf{1}, \frac{2}{3}\right) \oplus \left(\mathbf{1}, -\frac{2}{3}\right)$$

where the charges include the \mathcal{J} correction

- all gravitini carry *fractional electric charges!*
- complex triplet of gravitini is in addition strongly interacting
- the assignment makes the lightest of them stable – there is no particle it can decay into

Massive Gravitini

K.A.M., H. Nicolai, Phys. Rev. Lett. 121 (2018) 091601

- The eight massive gravitini under $SU(3) \times U(1)_{em}$ as

$$\left(\mathbf{3}, \frac{1}{3}\right) \oplus \left(\bar{\mathbf{3}}, -\frac{1}{3}\right) \oplus \left(\mathbf{1}, \frac{2}{3}\right) \oplus \left(\mathbf{1}, -\frac{2}{3}\right)$$

where the charges include the \mathcal{J} correction

- all gravitini carry *fractional electric charges!*
- complex triplet of gravitini is in addition strongly interacting
- the assignment makes the lightest of them stable – there is no particle it can decay into
- being stable, the color neutral gravitini should be around us (they were never in thermal equilibrium) – DM candidates, seeds of massive BHs in the very early Universe

Massive Gravitini

K.A.M., H. Nicolai, Phys. Rev. Lett. 121 (2018) 091601

- The eight massive gravitini under $SU(3) \times U(1)_{em}$ as

$$\left(\mathbf{3}, \frac{1}{3}\right) \oplus \left(\bar{\mathbf{3}}, -\frac{1}{3}\right) \oplus \left(\mathbf{1}, \frac{2}{3}\right) \oplus \left(\mathbf{1}, -\frac{2}{3}\right)$$

where the charges include the \mathcal{J} correction

- all gravitini carry *fractional electric charges!*
- complex triplet of gravitini is in addition strongly interacting
- the assignment makes the lightest of them stable – there is no particle it can decay into
- being stable, the color neutral gravitini should be around us (they were never in thermal equilibrium) – DM candidates, seeds of massive BHs in the very early Universe
- strongly interacting gravitini ('gravimesons') should also be around us but in much lower abundance – UHECR

Gravitini as DM

K.A.M. and H. Nicolai, Phys. Rev. D100 (2019) 035001

- Electrically charged DM are very strongly constrained by existing data:

$$|q| \lesssim 7.6 \cdot 10^{-10} \left(\frac{m}{1\text{TeV}} \right)^{\frac{1}{2}}$$

Gravitini as DM

K.A.M. and H. Nicolai, Phys. Rev. D100 (2019) 035001

- Electrically charged DM are very strongly constrained by existing data:

$$|q| \lesssim 7.6 \cdot 10^{-10} \left(\frac{m}{1\text{TeV}} \right)^{\frac{1}{2}}$$

- For the DM candidates usually discussed (axion-like or WIMP-like) assumed to have masses $\lesssim O(1)$ TeV the allowed charges are extremely small.

Gravitini as DM

K.A.M. and H. Nicolai, Phys. Rev. D100 (2019) 035001

- Electrically charged DM are very strongly constrained by existing data:

$$|q| \lesssim 7.6 \cdot 10^{-10} \left(\frac{m}{1\text{TeV}} \right)^{\frac{1}{2}}$$

- For the DM candidates usually discussed (axion-like or WIMP-like) assumed to have masses $\lesssim O(1)$ TeV the allowed charges are extremely small.
- If we extrapolate this formula to the Planck scale then $|q| \lesssim 1$ i.e. compatible with charges of our gravitini $\pm 2/3$ so our DM can be charged!

Gravitini as DM

K.A.M. and H. Nicolai, Phys. Rev. D100 (2019) 035001

- Electrically charged DM are very strongly constrained by existing data:

$$|q| \lesssim 7.6 \cdot 10^{-10} \left(\frac{m}{1\text{TeV}} \right)^{\frac{1}{2}}$$

- For the DM candidates usually discussed (axion-like or WIMP-like) assumed to have masses $\lesssim O(1)$ TeV the allowed charges are extremely small.
- If we extrapolate this formula to the Planck scale then $|q| \lesssim 1$ i.e. compatible with charges of our gravitini $\pm 2/3$ so our DM can be charged!
- these DM gravitini, being charged and very heavy, have very distinctive features

Search for DM

- The first argument of Zwicky from 30' still valid

Search for DM

- The first argument of Zwicky from 30' still valid
- LIGO/Virgo result combined with the other observations – 32% of matter: 5% luminous, 27% dark

Search for DM

- The first argument of Zwicky from 30' still valid
- LIGO/Virgo result combined with the other observations – 32% of matter: 5% luminous, 27% dark
- Thousands of theoretical papers, billions of dollars for experiments – zero result what DM is made of

Search for DM

- The first argument of Zwicky from 30' still valid
- LIGO/Virgo result combined with the other observations – 32% of matter: 5% luminous, 27% dark
- Thousands of theoretical papers, billions of dollars for experiments – zero result what DM is made of
- Allowed window for masses: 40 orders of magnitude 10^{-12} eV... 10^{28} eV

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.
- we assume that the gravitini are of \sim Planck mass

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.
- we assume that the gravitini are of \sim Planck mass
- if velocity $\sim 400 \text{ km/s}$ we arrive at a flux estimate

$$\Phi \sim 10^{-9} \text{ m}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \sim 0.03 \text{ m}^{-2} \text{ yr}^{-1} \text{ sr}^{-1}$$

(may be lower if DM co-rotates with the Solar system)

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.
- we assume that the gravitini are of \sim Planck mass
- if velocity $\sim 400 \text{ km/s}$ we arrive at a flux estimate

$$\Phi \sim 10^{-9} \text{ m}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \sim 0.03 \text{ m}^{-2} \text{ yr}^{-1} \text{ sr}^{-1}$$

(may be lower if DM co-rotates with the Solar system)

- large detectors (CMS, ATLAS or Superkamiokande).have the triggering focused only on relativistic particles

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.
- we assume that the gravitini are of \sim Planck mass
- if velocity $\sim 400 \text{ km/s}$ we arrive at a flux estimate

$$\Phi \sim 10^{-9} \text{ m}^{-2}\text{s}^{-1}\text{sr}^{-1} \sim 0.03 \text{ m}^{-2}\text{yr}^{-1}\text{sr}^{-1}$$

(may be lower if DM co-rotates with the Solar system)

- large detectors (CMS, ATLAS or Superkamiokande).have the triggering focused only on relativistic particles
- A dedicated time-of-flight experiment would be the best but difficult because of low flux

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.
- we assume that the gravitini are of \sim Planck mass
- if velocity $\sim 400 \text{ km/s}$ we arrive at a flux estimate

$$\Phi \sim 10^{-9} \text{ m}^{-2}\text{s}^{-1}\text{sr}^{-1} \sim 0.03 \text{ m}^{-2}\text{yr}^{-1}\text{sr}^{-1}$$

(may be lower if DM co-rotates with the Solar system)

- large detectors (CMS, ATLAS or Superkamiokande).have the triggering focused only on relativistic particles
- A dedicated time-of-flight experiment would be the best but difficult because of low flux
- The other more promising are 'paleodetectors' – looking for ionizing tracks in old crystals

Prospects of detection

- estimated mass density of DM in the proximity of the Solar System is $\sim 0.3 \cdot 10^6 \text{ GeV/m}^{-3}$.
- we assume that the gravitini are of \sim Planck mass
- if velocity $\sim 400 \text{ km/s}$ we arrive at a flux estimate

$$\Phi \sim 10^{-9} \text{ m}^{-2}\text{s}^{-1}\text{sr}^{-1} \sim 0.03 \text{ m}^{-2}\text{yr}^{-1}\text{sr}^{-1}$$

(may be lower if DM co-rotates with the Solar system)

- large detectors (CMS, ATLAS or Superkamiokande).have the triggering focused only on relativistic particles
- A dedicated time-of-flight experiment would be the best but difficult because of low flux
- The other more promising are 'paleodetectors' – looking for ionizing tracks in old crystals
- signature should be very different from anything else known

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- The gravitini in ordinary stars essentially do not annihilate, the cross section is too small

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- The gravitini in ordinary stars essentially do not annihilate, the cross section is too small
- However, if a star collapses to a neutron star, then the annihilation starts

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- The gravitini in ordinary stars essentially do not annihilate, the cross section is too small
- However, if a star collapses to a neutron star, then the annihilation starts
- The products can escape only when the annihilation takes place in the 'skin' of the neutron star (crust \sim last 100 m)

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- The gravitini in ordinary stars essentially do not annihilate, the cross section is too small
- However, if a star collapses to a neutron star, then the annihilation starts
- The products can escape only when the annihilation takes place in the 'skin' of the neutron star (crust \sim last 100 m)
- But in the crust there are mostly iron nuclei – the products of the collisions should be light nuclei and not protons!

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- The gravitini in ordinary stars essentially do not annihilate, the cross section is too small
- However, if a star collapses to a neutron star, then the annihilation starts
- The products can escape only when the annihilation takes place in the 'skin' of the neutron star (crust \sim last 100 m)
- But in the crust there are mostly iron nuclei – the products of the collisions should be light nuclei and not protons!
- extrapolating the formula for multiplicities from the LHC

$$\text{multiplicity} \sim 0.27 \alpha_s(E) \exp\left(\frac{2.26}{\sqrt{\alpha_s(E)}}\right)$$

to E_P we get $\sim 10^6$ (\Rightarrow particle energies 10^{21} eV)

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- It is exactly what is measured in Pierre Auger observatory UHECR ($\sim 10^{21}$ eV) are not protons but nuclei up to iron!

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- It is exactly what is measured in Pierre Auger observatory UHECR ($\sim 10^{21}$ eV) are not protons but nuclei up to iron!
- we calculate the flux:
 - For strongly interacting particles the annihilation cross section σ varies only very slowly with the energy \sqrt{s} , and can be approximated by the non-perturbative (Froissart bound) formula

$$\langle\sigma\beta\rangle\sim\left[36-4\ln\left(\frac{\sqrt{s}}{\Lambda_{QCD}}\right)+0.84\left(\ln\left(\frac{\sqrt{s}}{\Lambda_{QCD}}\right)\right)^2\right]\text{mb}$$

with $\Lambda_{QCD} = 0.4$ GeV. For $\sqrt{s} = 2M_P$ we have
 $\langle\sigma\beta\rangle \sim 32\text{mb}$

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- the relic abundance of color gravitini ρ_T at freeze-out

$$(32 \text{ mb}) \rho_T \equiv (32 \text{ mb}) g \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T} = \frac{T^2}{2M_P}$$

\Rightarrow

$$\frac{m}{T} \sim 90 \quad \Rightarrow \quad \rho_T \sim 3 \cdot 10^{59} \text{ m}^{-3}$$

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- the relic abundance of color gravitini ρ_T at freeze-out

$$(32 \text{ mb}) \rho_T \equiv (32 \text{ mb}) g \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T} = \frac{T^2}{2M_P}$$

\Rightarrow

$$\frac{m}{T} \sim 90 \quad \Rightarrow \quad \rho_T \sim 3 \cdot 10^{59} \text{ m}^{-3}$$

- The temperature $T \sim 2 \cdot 10^{16}$ GeV corresponds to cosmic time $t_T = M_P/T^2 \sim 3 \cdot 10^{-39}$ s.

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- the relic abundance of color gravitini ρ_T at freeze-out

$$(32 \text{ mb}) \rho_T \equiv (32 \text{ mb}) g \left(\frac{mT}{2\pi} \right)^{3/2} e^{-m/T} = \frac{T^2}{2M_P}$$

\Rightarrow

$$\frac{m}{T} \sim 90 \quad \Rightarrow \quad \rho_T \sim 3 \cdot 10^{59} \text{ m}^{-3}$$

- The temperature $T \sim 2 \cdot 10^{16}$ GeV corresponds to cosmic time $t_T = M_P/T^2 \sim 3 \cdot 10^{-39}$ s.
- the present day density of strongly interacting gravitini $\rho_T(a_T)/a_0^3 \sim 10^{-9} \text{ m}^{-3}$

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- density of gravitini inside a neutron star

$$\rho_{NS} \sim 5 \cdot 10^9 \text{ m}^{-3}.$$

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- density of gravitini inside a neutron star

$$\rho_{NS} \sim 5 \cdot 10^9 \text{ m}^{-3}.$$

- The inverse lifetime of the strongly interacting gravitino as a function of the neutron star time from its birth is

$$\Gamma_{NS}(t) = \frac{\Gamma_{NS}(0)}{1 + \Gamma_{NS}(0)t}$$

with the initial value (and $\langle \sigma \beta \rangle \sim 32 \text{ mb}$)

$$\Gamma_{NS}(0) \sim (5 \cdot 10^9) \cdot (32 \cdot 10^{-31}) \cdot (3 \cdot 10^8) \text{ s}^{-1} \sim 5 \cdot 10^{-12} \text{ s}^{-1}$$

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- number of neutron stars per galaxy $\sim 10^8$;
 $\sim 10^9$ galaxies within ~ 250 Mpc;
a total number of 10^{17} UHECR emitters.

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- number of neutron stars per galaxy $\sim 10^8$;
 $\sim 10^9$ galaxies within ~ 250 Mpc;
a total number of 10^{17} UHECR emitters.
- a young neutron star would continuously ‘spray’ high energy protons or heavy ions at a rate $\sim 10^{16} \text{s}^{-1}$

Ultra High Energy Cosmic Rays

K.A.M. and H. Nicolai, JCAP 09 (2019) 041

- number of neutron stars per galaxy $\sim 10^8$;
 $\sim 10^9$ galaxies within ~ 250 Mpc;
a total number of 10^{17} UHECR emitters.
- a young neutron star would continuously ‘spray’ high energy protons or heavy ions at a rate $\sim 10^{16}\text{s}^{-1}$
- we estimate the flux

$$N_E \sim \frac{10^{17} \cdot 10^{16}}{4(10^{24})^2} \text{ m}^{-2}\text{s}^{-1} \sim 10^{-16} \text{ m}^{-2}\text{s}^{-1}$$

which is close to the observed rate of one UHECR event per month and per 3000 km^2

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- Very massive (> 1 bln M_{\odot}) black holes are observed in the very early (< 1 bln y) Universe.

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- Very massive ($> 1 \text{ bln } M_{\odot}$) black holes are observed in the very early ($< 1 \text{ bln y}$) Universe.
- no standard mechanism can explain the observations

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- Very massive ($> 1 \text{ bln } M_{\odot}$) black holes are observed in the very early ($< 1 \text{ bln y}$) Universe.
- no standard mechanism can explain the observations
- (Color neutral) gravitini can form an initial black hole large enough to be colder than the surrounding and overcome Hawking evaporation

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- Very massive ($> 1 \text{ bln } M_{\odot}$) black holes are observed in the very early ($< 1 \text{ bln y}$) Universe.
- no standard mechanism can explain the observations
- (Color neutral) gravitini can form an initial black hole large enough to be colder than the surrounding and overcome Hawking evaporation
- Using the solution

$$ds^2 = a(\eta)^2 \left[-(1 - 2m_{\text{BH}}/r)d\eta^2 + \frac{dr^2}{1 - 2m_{\text{BH}}/r} + r^2 d\Omega^2 \right]$$

in the expanding Universe we get the bounds at $t \sim 100 \text{ Myr}$

$$10^5 M_{\odot} \lesssim m_{\text{BH}} \lesssim 2 \cdot 10^9 M_{\odot}$$

which is consistent with observations

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- As a consequence of the Hawking evaporation of seed black holes with too small mass, our calculation also provides a *lower bound*

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- As a consequence of the Hawking evaporation of seed black holes with too small mass, our calculation also provides a *lower bound*
- It is thus a prediction of the present mechanism that the black holes formed from gravitinos should belong to a very different mass category than the black holes formed from stellar collapse and subsequent mergers

Gravitini and giant black holes

K.A.M., H. Nicolai, Phys.Rev. D102 (2020) 103008

- As a consequence of the Hawking evaporation of seed black holes with too small mass, our calculation also provides a *lower bound*
- It is thus a prediction of the present mechanism that the black holes formed from gravitinos should belong to a very different mass category than the black holes formed from stellar collapse and subsequent mergers
- such a gap in the mass distribution of black holes in the Universe would constitute indirect observational evidence for the existence of Hawking radiation.

Summary

- adjustment of $N = 8 \rightarrow SU(3) \times U(1)$ that gives proper assignment of electric charges to quarks and leptons requires that 8 very heavy gravitini are fractionally charged, stable, strongly and electromagnetically interacting

Summary

- adjustment of $N = 8 \rightarrow SU(3) \times U(1)$ that gives proper assignment of electric charges to quarks and leptons requires that 8 very heavy gravitini are fractionally charged, stable, strongly and electromagnetically interacting
- being extremely massive (and therefore very diluted) but stable they can be candidates for (not so) Dark Matter

Summary

- adjustment of $N = 8 \rightarrow SU(3) \times U(1)$ that gives proper assignment of electric charges to quarks and leptons requires that 8 very heavy gravitini are fractionally charged, stable, strongly and electromagnetically interacting
- being extremely massive (and therefore very diluted) but stable they can be candidates for (not so) Dark Matter
- their annihilation in neutron stars can explain the origin of UHECR (in the form of light nuclei) observed on Earth

Summary

- adjustment of $N = 8 \rightarrow SU(3) \times U(1)$ that gives proper assignment of electric charges to quarks and leptons requires that 8 very heavy gravitini are fractionally charged, stable, strongly and electromagnetically interacting
- being extremely massive (and therefore very diluted) but stable they can be candidates for (not so) Dark Matter
- their annihilation in neutron stars can explain the origin of UHECR (in the form of light nuclei) observed on Earth
- Their clumping in the very early Universe can also explain the origin of observed giant black holes in the early Universe

Summary

- adjustment of $N = 8 \rightarrow SU(3) \times U(1)$ that gives proper assignment of electric charges to quarks and leptons requires that 8 very heavy gravitini are fractionally charged, stable, strongly and electromagnetically interacting
- being extremely massive (and therefore very diluted) but stable they can be candidates for (not so) Dark Matter
- their annihilation in neutron stars can explain the origin of UHECR (in the form of light nuclei) observed on Earth
- Their clumping in the very early Universe can also explain the origin of observed giant black holes in the early Universe
- time will tell...

Happy Birthday, Hermann!