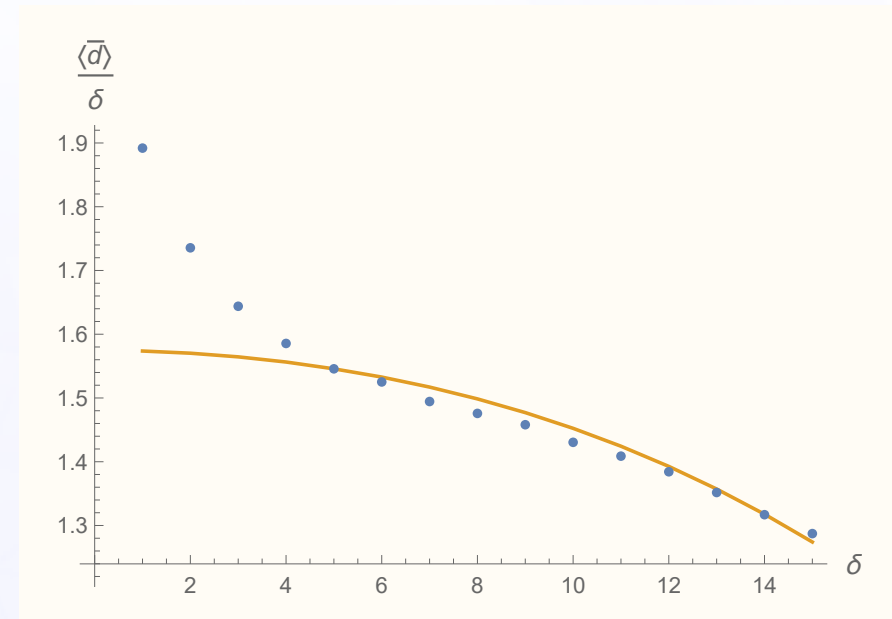


J. Ambjørn, J. Jurkiewicz, R.L., PRL 93 (2004) 131301



N. Klitgaard, R.L., EPJ C 80 (2020) 990

Quantum Gravity Demystified

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My perspective on quantum gravity

Aim: construct a fundamental theory of quantum gravity as a non-perturbative, diffeomorphism-invariant quantum field theory of dynamical geometry and study its properties in a Planckian regime.

This presents major technical, physical and conceptual challenges: dealing with QFT infinities and the absence of a fixed background spacetime, devising appropriate numerical and renormalization methods, (re-)deriving the classical limit and phenomenology.

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This is possible. Major advances towards this goal have been made in the research program of ***Causal Dynamical Triangulations (CDT)***. It sets a concrete frame of reference - beyond matters of taste and style - for what we may reasonably expect to be able to achieve in quantum gravity.

Why should you care?

quantum gravity: nontrivial, unexpected results despite non-exotic ingredients; functioning computational framework (= *our “lab”*) to evaluate quantum observables beyond perturbation theory;

“CDT is to gravity what lattice QCD is to nonabelian gauge theory”

symmetry: diffeomorphism symmetry is very different from local gauge symmetry; we finally understand how to implement it consistently in a nonperturbative quantum theory of gravity

“demystification”: quantum (field) theory and general relativity are perfectly compatible; CDT provides bottom-up realization of QG: causal structure is essential, unitarity is realized

cosmology: most likely phenomenological predictions will involve early-universe quantum physics, but derived from the full theory *without* an a priori symmetry reduction (unlike quantum cosmology)

What's the problem with quantum gravity?

- General Relativity = theory **of** spacetime, not **on** (a fixed) spacetime
- quantum theory based on perturbative split $g_{\mu\nu}(x) = \eta_{\mu\nu}^{\text{Mink}} + h_{\mu\nu}(x)$ on a fixed Minkowskian background is nonrenormalizable

M. Goroff, A. Sagnotti, NPB 266 (1986) 709

- standard relativistic quantum field theory (QFT) not applicable, no blueprint beyond perturbation theory (except nonperturbative lattice QCD, **but** this has a fixed background, different gauge symmetry)
- no experiments or observations to guide theory-building
- (nonperturbative) QG \approx 2000: large variety of approaches;
“**We don't know *what* to compute, and we don't know *how*.**”
- QG \approx 2000 (post extended-objects era): renaissance of “good old QFT”/the path integral, *we have learned how and what to compute*

R.L. et al.: “Quantum Gravity in 30 Questions”, arXiv: 2206.06762

Causal Dynamical Triangulations: the basics

- gravitational path integral over metric d.o.f., *nonperturbative* (NP), background-independent, Lorentzian signature, 4D, not “grand-unified”

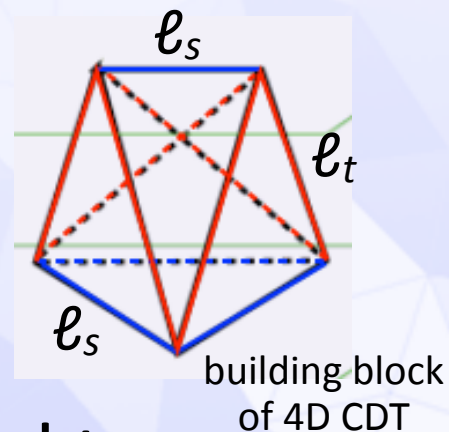
$$Z = \int_{g \in \frac{Lor(M)}{Diff(M)}} \mathcal{D}g \, e^{iS_{\text{grav}}[g]}$$

- building on Euclidean “dynamical triangulations”, define a new NP Lorentzian 2D path integral: *exactly soluble* \Rightarrow *signature matters!*

J. Ambjørn, R.L., preprint AEI-064, NPB 536 (1998) 407

- CDT combines emphasis on geometry with path integral covariance (no split $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$, no 3 + 1 decomposition)

- uses a regularized version of the space of geometries, $\mathcal{G}(M) = Lor(M)/Diff(M)$: piecewise flat, simplicial manifolds \mathcal{T}



- minimal GR ingredients + standard Q(F)T methods, adapted to dynamical geometry + numerical methods = **new territory near ℓ_{PI}**

- 2D random geometry is a hot topic in maths! S. Sheffield, arXiv:2203.02470

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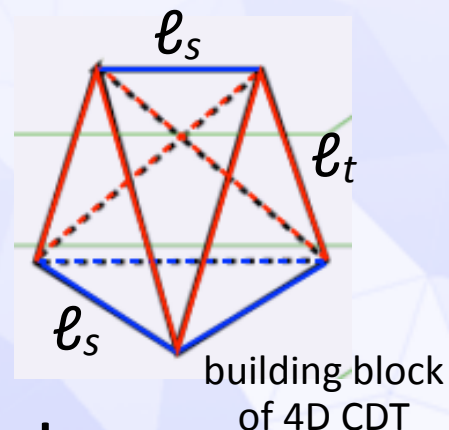
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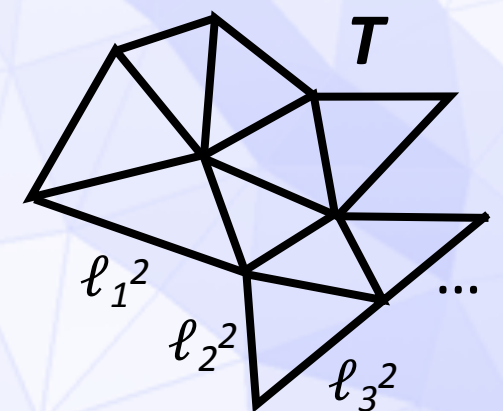
Putting quantum gravity on a lattice, correctly

General strategy: lattice acts as a regulator, with UV cutoff a ; search for a continuum limit by approaching a second-order phase transition in the limit $a \rightarrow 0$ while renormalizing bare couplings appropriately; attain *universality* (independence of regularization); this is **not** “discrete QG”

- “reaches where other methods don’t”, subject to numerical limitations; if it exists, continuum theory is essentially *unique*
- “naïve” lattice QG (≥ 1979): put various first-order formulations of GR (tetrad e_μ^A + spin connection ω_μ^{AB}) on a fixed hypercubic lattice; problem: diffeomorphism symmetry badly broken; no interesting results
- “not-so-naïve” lattice QG (≥ 1981): based on “GR without coordinates”
 $(M, g_{\mu\nu}(x)) \rightarrow (T, \{\ell_i^2, i=1, \dots, n\}), S_{\text{grav}}[g_{\mu\nu}] \rightarrow S^{\text{Regge}}(T, \{\ell_i^2\})$

T. Regge, *Nuovo Cim.* A19 (1961) 558

- diffeo-invariance manifest, work directly on $\mathcal{G}(M)$;
CDT ($\ell^2 = \pm a^2$) implementation is labelling-invariant



The path integral (PI) according to CDT

$$Z = \int_{\mathcal{G}(M)} \mathcal{D}g \, e^{iS_{\text{grav}}[g]} \rightarrow Z^{\text{CDT}} = \lim_{a \rightarrow 0} \sum_{\substack{\text{inequiv.} \\ \text{causal} \\ \text{triang. } T}} \frac{1}{C(T)} e^{iS^{\text{Regge}}[T]}$$

bare action
discrete symmetries of T

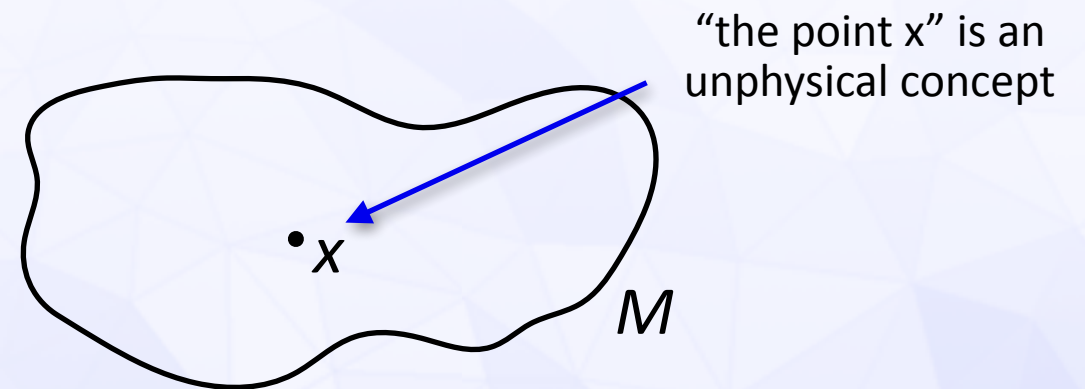
- usually, can't evaluate complex PI, do Euclidean $\int Dg \exp(-S^{eu})$ instead
 - ☑ CDT has a well-defined analytic continuation ("Wick-rotation")
- usually, hard to renormalize compatible with diffeomorphism symmetry
 - ☑ CDT has no residual symmetries, has a geometric cutoff a
- usually, PI highly divergent, no unique renormalization;
 - ☑ numerical evidence of exponential bound on # of configurations
- usually, cannot do any computations, PI not Gaussian
 - ☑ CDT amenable to Monte Carlo simulation; get quantitative results
- usual problem: why should PI lead to a unitary theory?
 - ☑ CDT reflection-positive w.r.t. discrete "proper time"

CDT quantum gravity: results

- we have a computational framework — what can we do with it?
- physics of **quantum spacetime** is captured by invariant **quantum observables** $\hat{\mathcal{O}}$:

$$\langle \hat{\mathcal{O}} \rangle = \frac{1}{Z} \int \mathcal{D}g \, \mathcal{O}[g] e^{-S_{\text{grav}}[g]}$$

- observables in Yang-Mills theory are local scalars, like $F^{\mu\nu}F_{\mu\nu}$, but observables in pure gravity are nonlocal integrals of scalars, like $\int_M d^4x \sqrt{g} R(x)$



- “expectation management”: your favourite (semi-)classical question will not have a Planckian implementation (this is a feature)
- quantum gravity signature: CDT predicts a reduction $4 \rightarrow 2$ of the **spectral dimension** @ ℓ_{Pl} , [J. Ambjørn, J. Jurkiewicz, R.L., PRL 95 \(2005\) 171301](#) repro-duced across approaches — universal in QG? [S. Carlip, CQG 34 \(2017\) 193001](#)

Key result: emergence of classicality from CDT

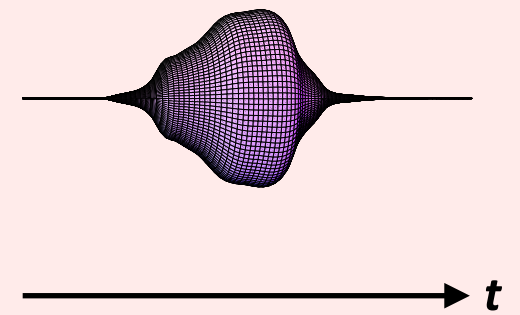
The measured average shape $\langle V_3(t) \rangle$ (spatial volume as a function of proper time) of the dynamically generated quantum spacetime in CDT matches that of a classical de Sitter space.

J. Ambjørn, A. Görlich, J. Jurkiewicz, R.L., PRL 100 (2008) 091304, PRD 78 (2008) 063544

Since the global shape of the universe is just a single mode of the metric, we cannot conclude that it *is* a (Euclidean) de Sitter space S^4 , with line element

$$ds^2 = dt^2 + c^2 \cos^2(t/c) d\Omega_{(3)}^2.$$

MC snapshot of the shape
 $\langle V(t) \rangle$ of the universe

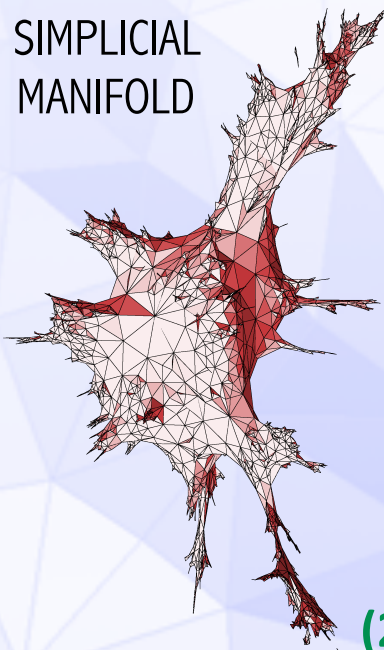


What about the **local geometry** of this quantum universe?

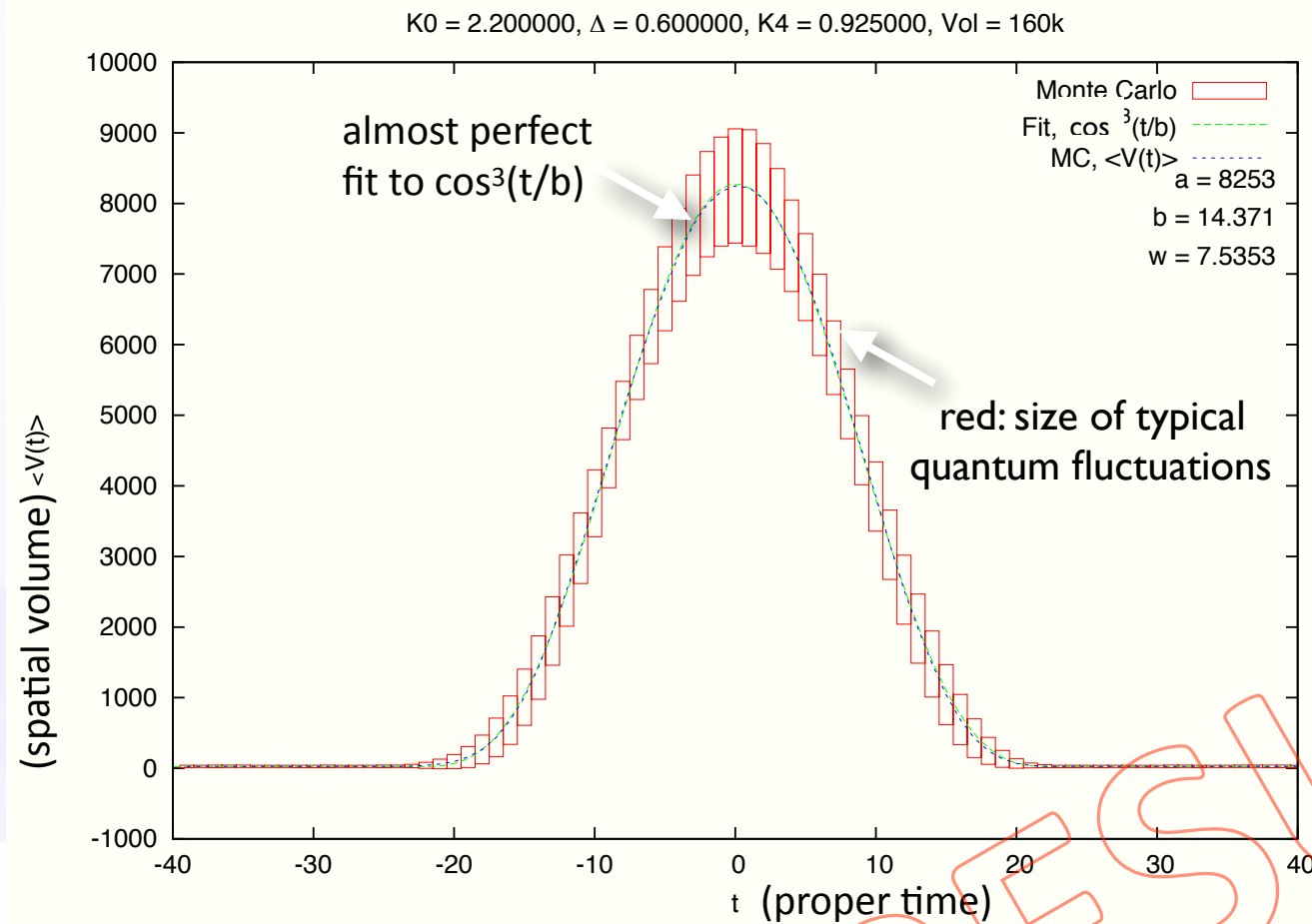
Can we attribute **local curvature** to a non-smooth metric space? $R^\kappa_{\lambda\mu\nu}(x) = [g, \partial g, \partial^2 g] = ?$ Recently, we defined, tested and measured a well-defined notion of **quantum Ricci curvature** applicable in a Planckian regime. N. Klitgaard & RL, PRD 97

(2018) 0460008 and 106017, Eur. Phys. J. C80 (2020) 990, J. Brunekreef & RL, PRD 103 (2021) 026019

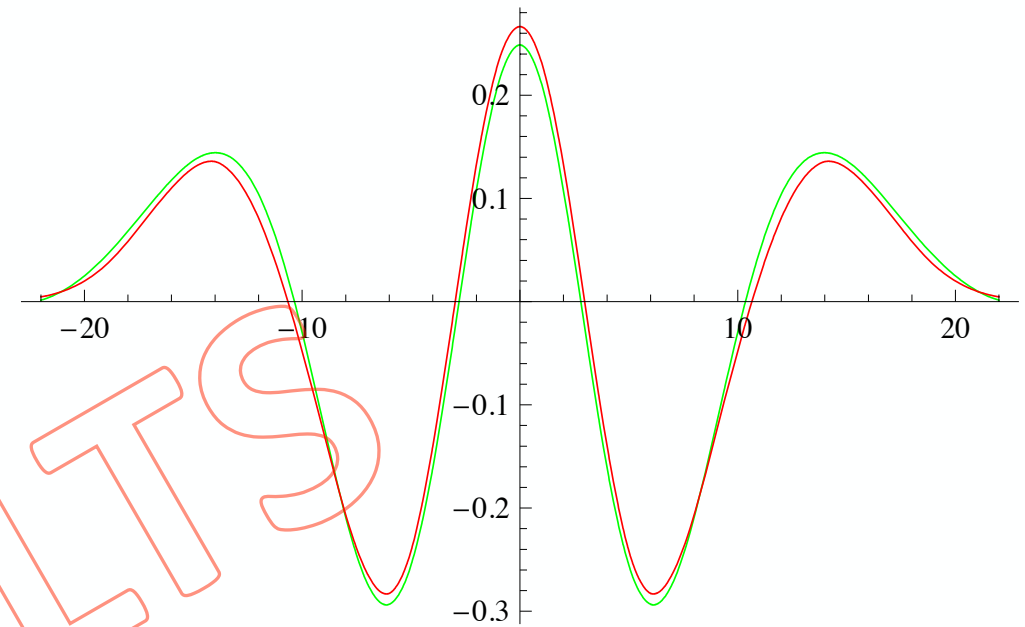
SIMPLICIAL
MANIFOLD



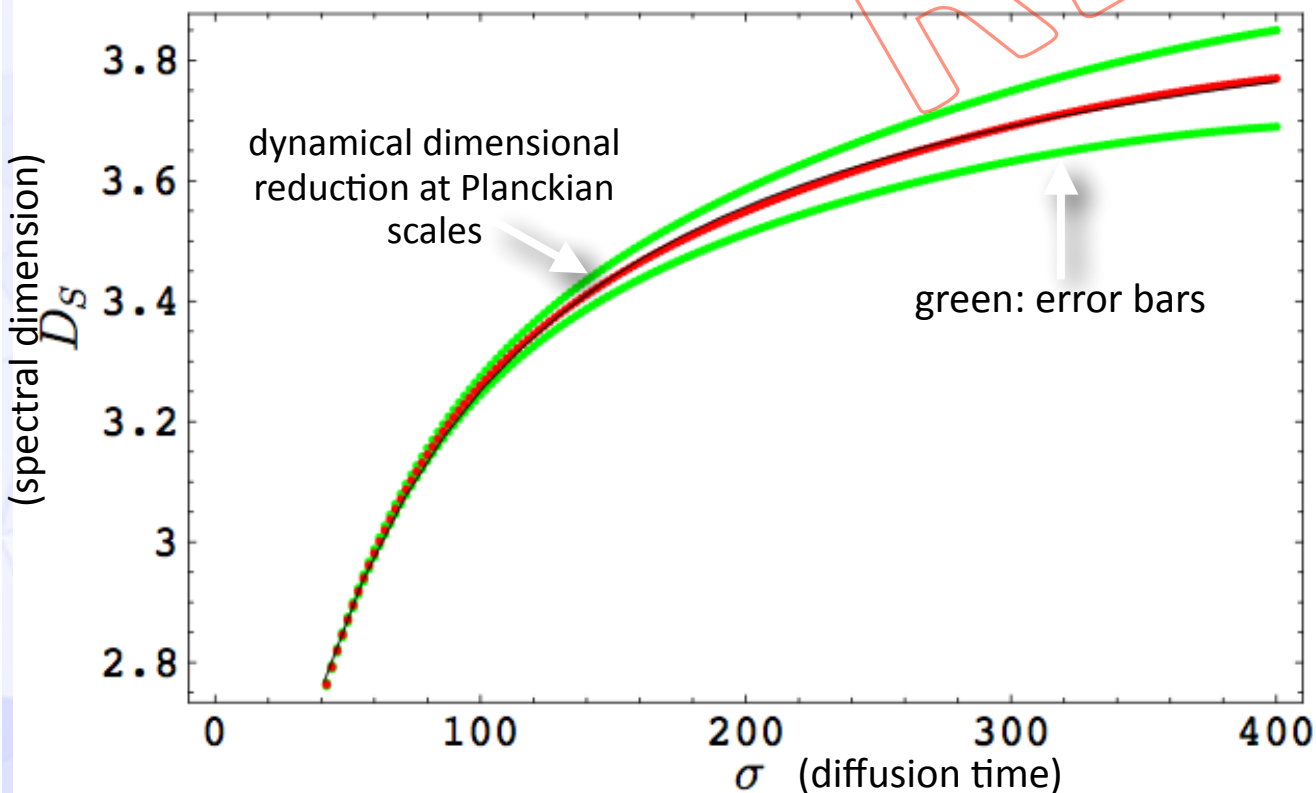
The universe is de Sitter-shaped



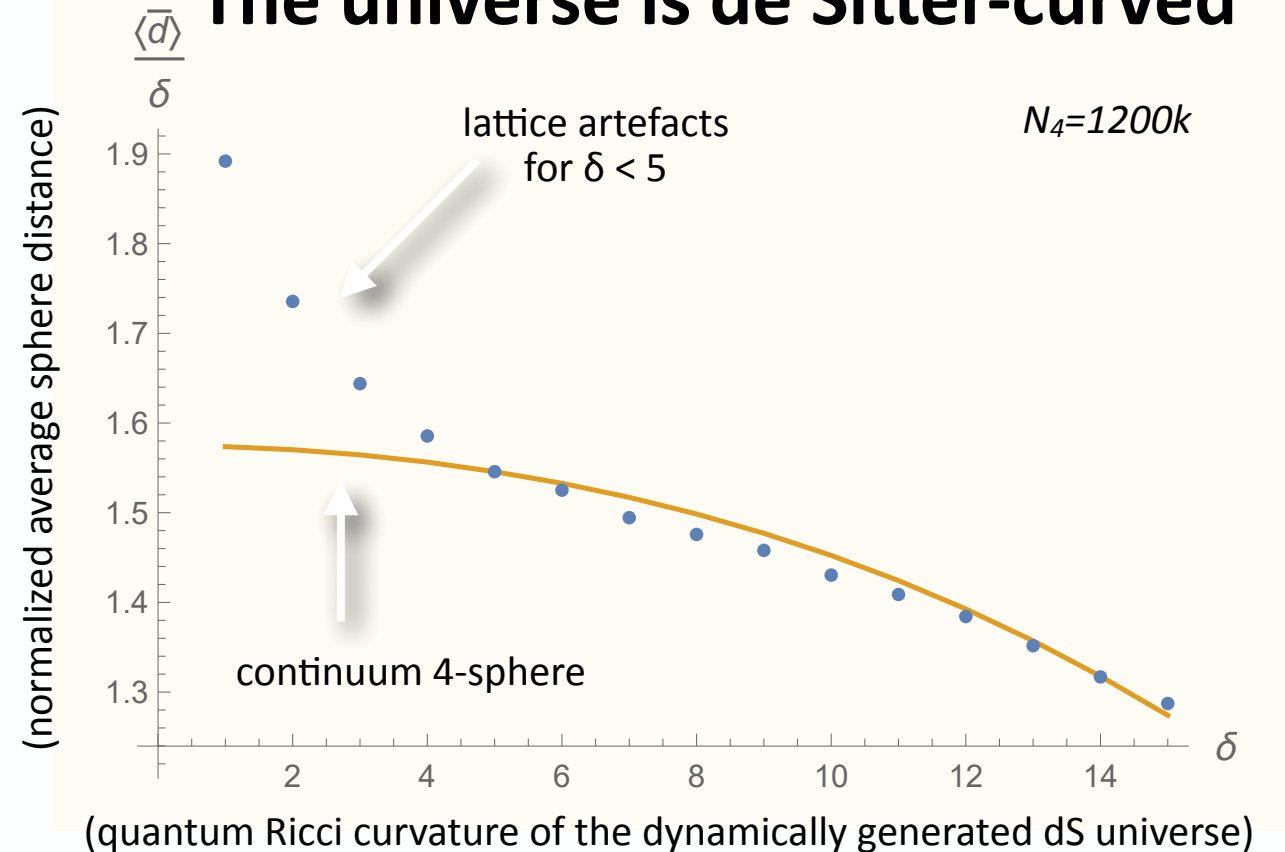
Volume fluctuations around de Sitter



Spectral dimension of the universe



The universe is de Sitter-curved



Relation to our actual universe

CDT predicts a universe **with positive Λ** , which on large scales is extended and four-dimensional, and whose **shape and average curvature** are compatible with those **of a de Sitter space**, matching our current understanding of the very early universe.

Remarkably, these properties have been derived *from first principles* in the full quantum theory; we also have in principle access to (diffeomorphism-invariant) correlation functions.

At what scales and how does gravity interact with matter?

$$Z = \int_{\mathcal{G}(M)} \mathcal{D}g \int_{\Phi} \mathcal{D}\phi \, e^{i(S_{\text{grav}}[g] + S_{\text{matter}}[g, \phi])}$$

Investigations of CDT coupled to matter fields have not found a significant impact on the geometry \Rightarrow “matter doesn’t matter at ℓ_{Pl} ”?

Current ambitions and prospects

CDT QG is in a position to reap the benefits of a nonperturbative framework that can produce “numbers” (= expectation values of quantum observables) without relying on ad-hoc assumptions. The art is to ***identify (more) observables*** that can be reliably measured inside the available scale window, while yielding interesting physics.

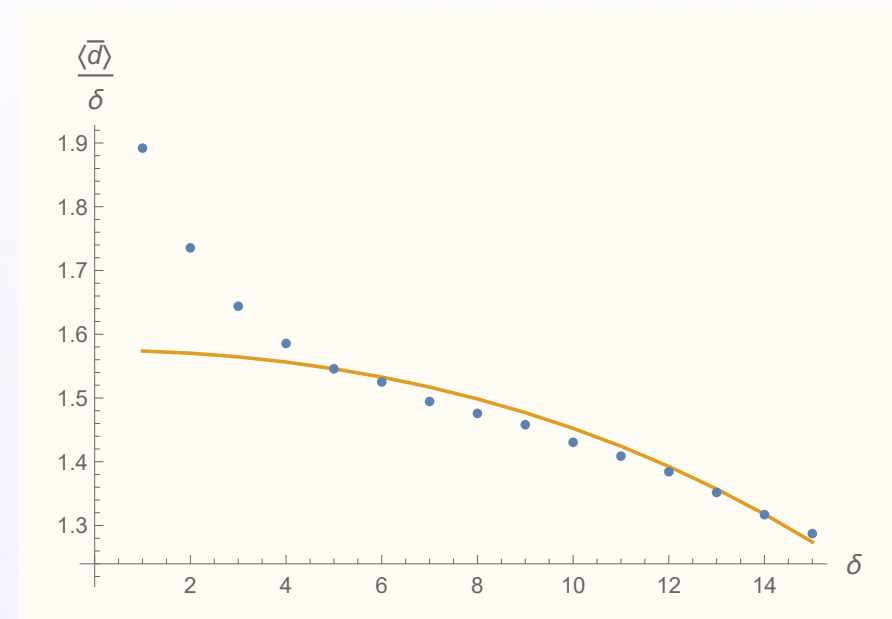
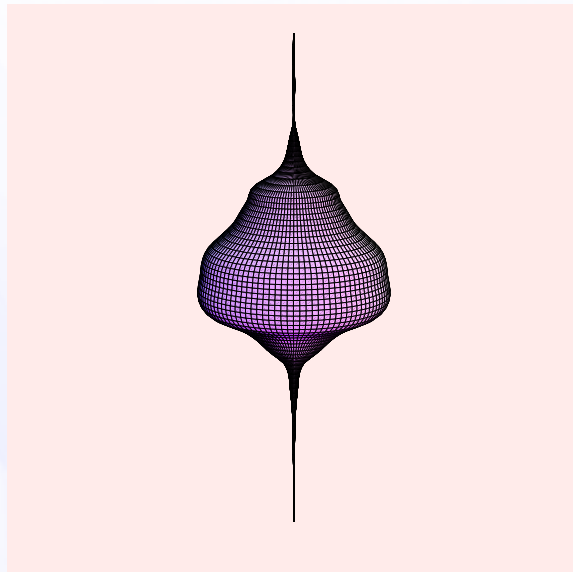
The new ***quantum Ricci curvature*** [Y. Ollivier, J. Funct. Anal. 256 \(2009\) 810](#) opens exciting avenues towards a relation with early-universe physics:

- use it to quantify the *local* effect of a massive Planckian particle and compare with a semiclassical expectation [G. Clemente, R.L., w.i.p.](#)
- use it to examine the string-like singularity spontaneously forming in the bifurcation phase of CDT, as a possible *candidate of early-universe structure formation (a primordial black hole?)*
- compute its *two-point function*, compare with QFT on dS space

Outlook

- ***genuine progress*** in nonperturbative quantum gravity: instead of comparing “approaches”, started to compare observables/results, e.g. with those obtained by functional RG methods [F. Saueressig et al.](#)
 - CDT quantum gravity is a rare example of ***spacetime emergence***
 - ***work in progress***: quantum measures of homogeneity & isotropy [A. Silva, R.L.](#); extend earlier RG flow analysis [J. Ambjørn, J. Gizbert-Studnicki, A. Görlich, J. Jurkiewicz, R.L., Front. in Phys. 8 \(2020\) 247](#) and look for independent evidence of asymptotic safety; measure curvature correlators, starting in 2D [J. van der Duin, R.L.](#)
 - ***challenge***: match nonperturbative and perturbative observables
- ⇒ watch this space!

CDT reviews: [J. Ambjørn, A. Görlich, J. Jurkiewicz, R.L., Phys. Rep. 519 \(2012\) 127, arXiv: 1203.3591](#); [R.L., Class. Quant. Grav. 37 \(2020\) 013002, arXiv:1905.08669](#)



Thank you!