

Extreme Mass Ratio Inspirals in non-Kerr spacetimes

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EMRIs in Kerr: a short summary

- No matter, Petrov type V spacetime
 - Geodesic equations separable
 - Efficient way to compute metric perturbations (Teukolsky formalism)
- Hard to compute inspiral phase:
 - Solve $\frac{Du^\mu}{d\tau} = -\frac{1}{2}(g^{\mu\nu} + u^\mu u^\nu)(2\nabla_\rho h_{\nu\lambda} - \nabla_\nu h_{\lambda\rho})u^\lambda u^\rho$:
metric perturbations diverging at the particle \Rightarrow
regularization needed \Rightarrow self-force approach: no inspirals
computed so far
 - Compute $\dot{E}, \dot{L}_z, \dot{Q}$: particle follows a geodesic with
evolving parameters.
Less general than self-force approach, problems with \dot{Q}
First waveforms in 2005

Going beyond Kerr: why?

- Exotic alternatives to Kerr SMBHs: gravastars, boson stars, fermion balls
- Even if the SMBH is actually a BH, there's a lot of matter in galactic centers: accretion disks in AGNs, stellar disks, dark matter density profile is peaked
- Little is known about the strong field region close to the horizon: can GWs from EMRIs be a tool to detect deviations from Kerr?

Must be prepared for surprises!

Our approach

- Numerical solutions of BH+Torus spacetimes now available (Ansorg & Petroff, Phys. Rev. D 72 024019)
 - Tori are heavy, compact and close to the horizon: unrealistic but useful as a strawman to maximize impact of “astrophysical bumpiness” on EMRIs
 - Approach follows in spirit Glampedakis and Babak's (CQG 23, 4167): compared kludge waveforms produced in Kerr and Kerr+quadrupole spacetimes; found that waveforms containing the same frequencies are indistinguishable by LISA (confusion problem)
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What we do: a summary

Numerical integration of geodesics
in BH+Torus spacetimes:

$$p = \frac{\text{latus rectum}}{r_{max} + r_{min}} = \frac{2 r_{max} r_{min}}{r_{max} + r_{min}} \quad e = \frac{\text{eccentricity}}{r_{max} + r_{min}} = \frac{r_{max} - r_{min}}{r_{max} + r_{min}} \quad z_- = \frac{\text{inclination}}{[\cos^2(\theta)]_{max}}$$

Compute kludge waveforms using geodesic motion
(with cutoff in time to avoid radiation reaction effects)
in quasi-BL coordinates

Address the confusion problem: compare BH+Torus
waveforms to the “corresponding” waveforms
produced in a Kerr spacetime

What are kludge waveforms?

Folk etymology: klumsy, lame, ugly, dumb, but good enough

- Identify BL or quasi BL-coordinates with spherical coordinates
 - Use $\partial_\nu T_{particle}^{\mu\nu} = 0$ (incorrect in curved spacetimes) + Green's function of \square in Minkowski spacetime (incorrect too)
 - Surprisingly good agreement with Teukolsky-based waveforms in Kerr
 - Natural preliminary approach in quasi-Kerr spacetimes, where Teukolsky formalism is not applicable
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How long can we trust our waveforms?

Consider an orbit with initial parameters p , e and z_-

Compute waveforms with (h_{RR})

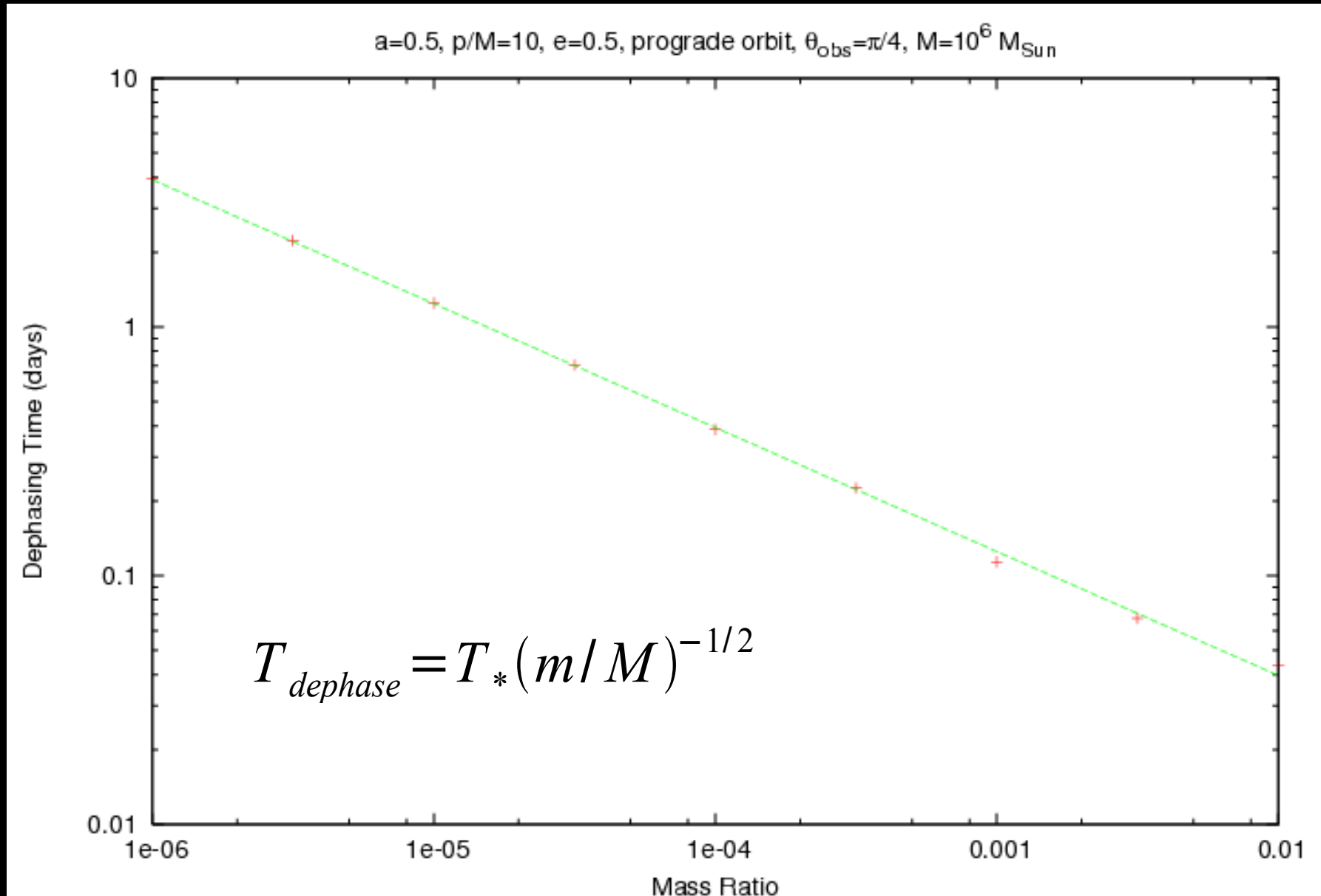
and without (h_{noRR}) radiation reaction:

radiation reaction becomes important when

$overlap(h_{noRR}, h_{RR})$ drops below 0.95 (dephasing time)

- In pure Kerr we can use “kludge” fluxes for E , L_z and Q (based on post-newtonian expansions and fits to Teukolsky data)
- In BH+Torus we use the dephasing time of the orbit in Kerr having the same p , e and z_- in quasi-Boyer-Lindquist coordinates

How long is the dephasing time?



Recalling the signal to noise ratio...

$$s(t) = h(t) + n(t)$$

signal waveform gaussian noise

$$\frac{S}{N}[\hat{h}] = \frac{\int \hat{h}(t) w(t - \tau) s(\tau) d\tau dt}{\text{rms} \int \hat{h}(t) w(t - \tau) n(\tau) d\tau dt} = \frac{(\hat{h}, s)}{(\hat{h}, \hat{h})^{1/2}}$$

template

$$\tilde{w}(f) = 1/S_n(f)$$

detector's sensitivity

$$(h_1, h_2) \equiv 2 \int_0^\infty \frac{\tilde{h}_1^*(f) \tilde{h}_2(f) + \tilde{h}_1(f) \tilde{h}_2^*(f)}{S_n(f)} df$$

The overlap function

Template matching the waveform ($\hat{h} = h$):

$$S/N = (h, h)^{1/2} \equiv A$$

Template not matching the waveform:

$$S/N = A \cdot \text{overlap}(h, \hat{h})$$

$$\text{overlap}(h_1, h_2) \equiv (h_1, h_2) / [(h_1, h_2)^{1/2} (h_1, h_2)^{1/2}]$$

Overlap between BH+Torus

and the “closest” Kerr waveform

expresses how much the S/N ratio is degraded

if we match-filter a BH+Torus signal with Kerr templates

Comparing Kerr and BH+Torus

- Compare waveforms having the same frequencies:
 - Get the frequencies for the BH+Torus orbit (“approximate” frequencies if non-equatorial)
 - Get the Kerr orbit having the same frequencies by solving the system

$$\Omega_r^{Kerr}(p, e, z_-, a, M) = \Omega_r^{BH+Torus}$$

$$\Omega_\varphi^{Kerr}(p, e, z_-, a, M) = \Omega_\varphi^{BH+Torus}$$

$$\Omega_\theta^{Kerr}(p, e, z_-, a, M) = \Omega_\theta^{BH+Torus}$$

in 3 variables chosen among p, e, z_-, a, M

(variables not solved for are fixed to their BH+Torus values, e. g. $p = p_{BH+Torus}$, $e = e_{BH+Torus}$, etc.)

An example: equatorial orbits

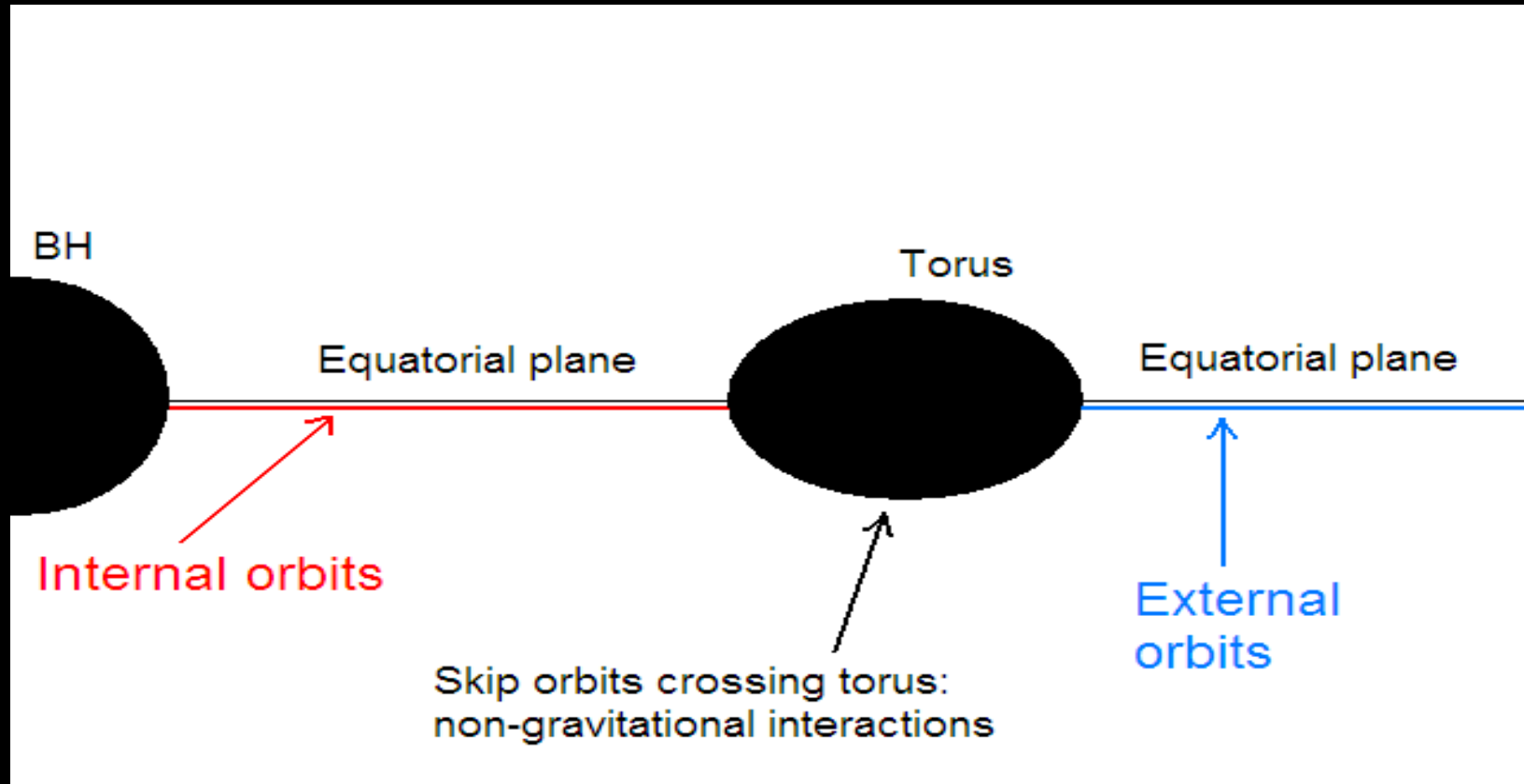
- 2 equations (equality of Ω_r and Ω_φ) and 4 variables (p, e, a, M)
 - Can solve for p and e and compute overlap between BH+Torus wfms and Kerr wfms with the same a, M, Ω_r and Ω_φ but different p and e :
if overlap > 0.95 wfms are indistinguishable
 (“confusion problem in p and e ”)
 - Can solve for a and M and compute overlap between BH+Torus wfms and Kerr wfms with the same p, e, Ω_r and Ω_φ but different a and M :
if overlap > 0.95 wfms are indistinguishable
 (“confusion problem in a and M ”)
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Spacetime A $M_{BH} = 10^6 M_{\odot}$ $M_{Torus} = 7 \cdot 10^4 M_{\odot}$
 $a_{BH} = -1.74 \cdot 10^{-3}$ $a_{BH+Torus} = 0.224$
 $r_{QI\ inner}/M_{BH} = 9.16$ $r_{QI\ outer}/M_{BH} = 10$
 $\epsilon = 2.63$

Spacetime B $M_{BH} = 4.13 \cdot 10^6 M_{\odot}$ $M_{Torus} = 1.21 \cdot 10^6 M_{\odot}$
 $a_{BH} = 0.528$ $a_{BH+Torus} = 0.728$
 $r_{QI\ inner}/M_{BH} = 1.468$ $r_{QI\ outer}/M_{BH} = 1.527$
 $\epsilon = 0.1103$

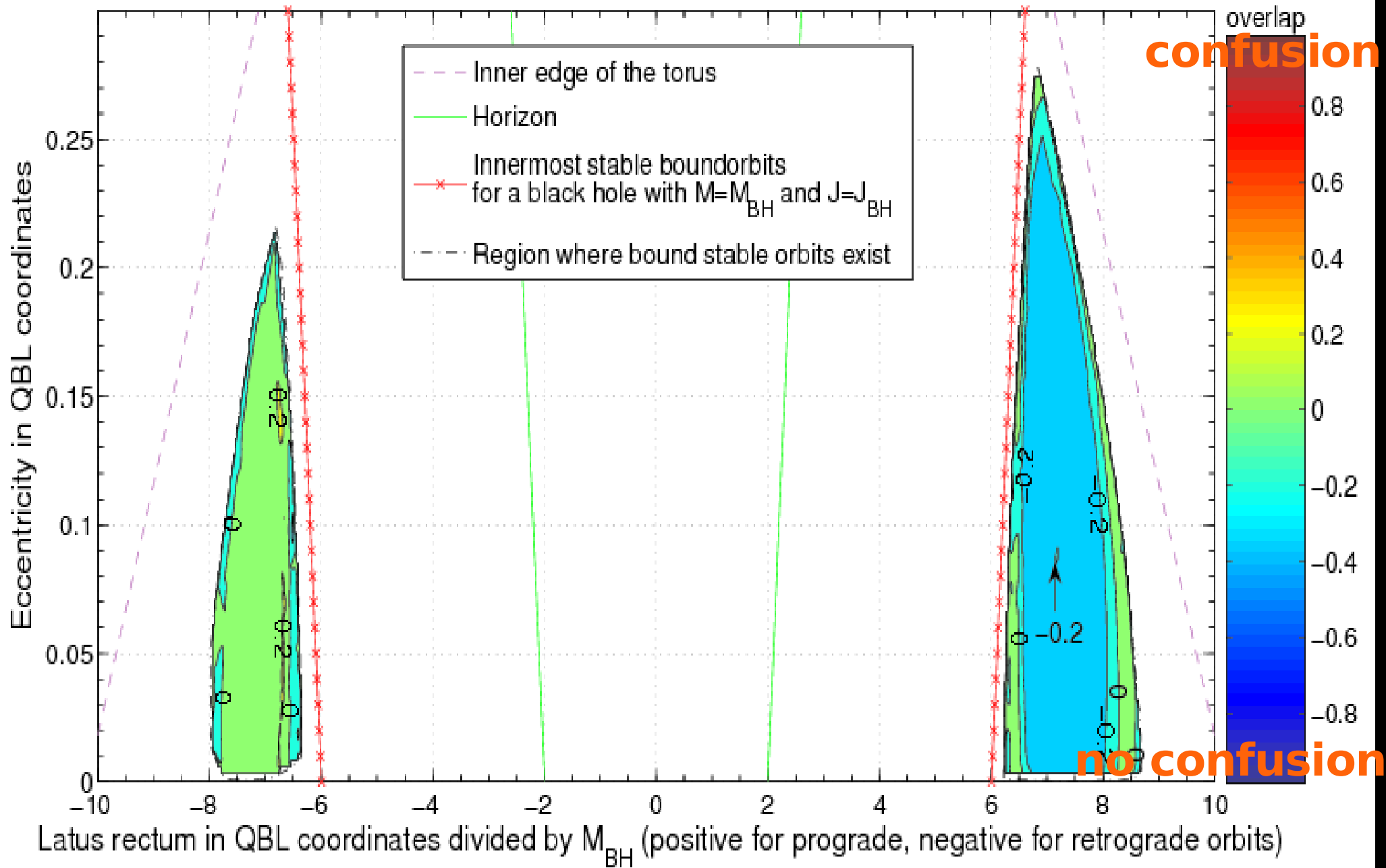
Small BH $m = 1 M_{\odot}$

Equatorial orbits: summary of results

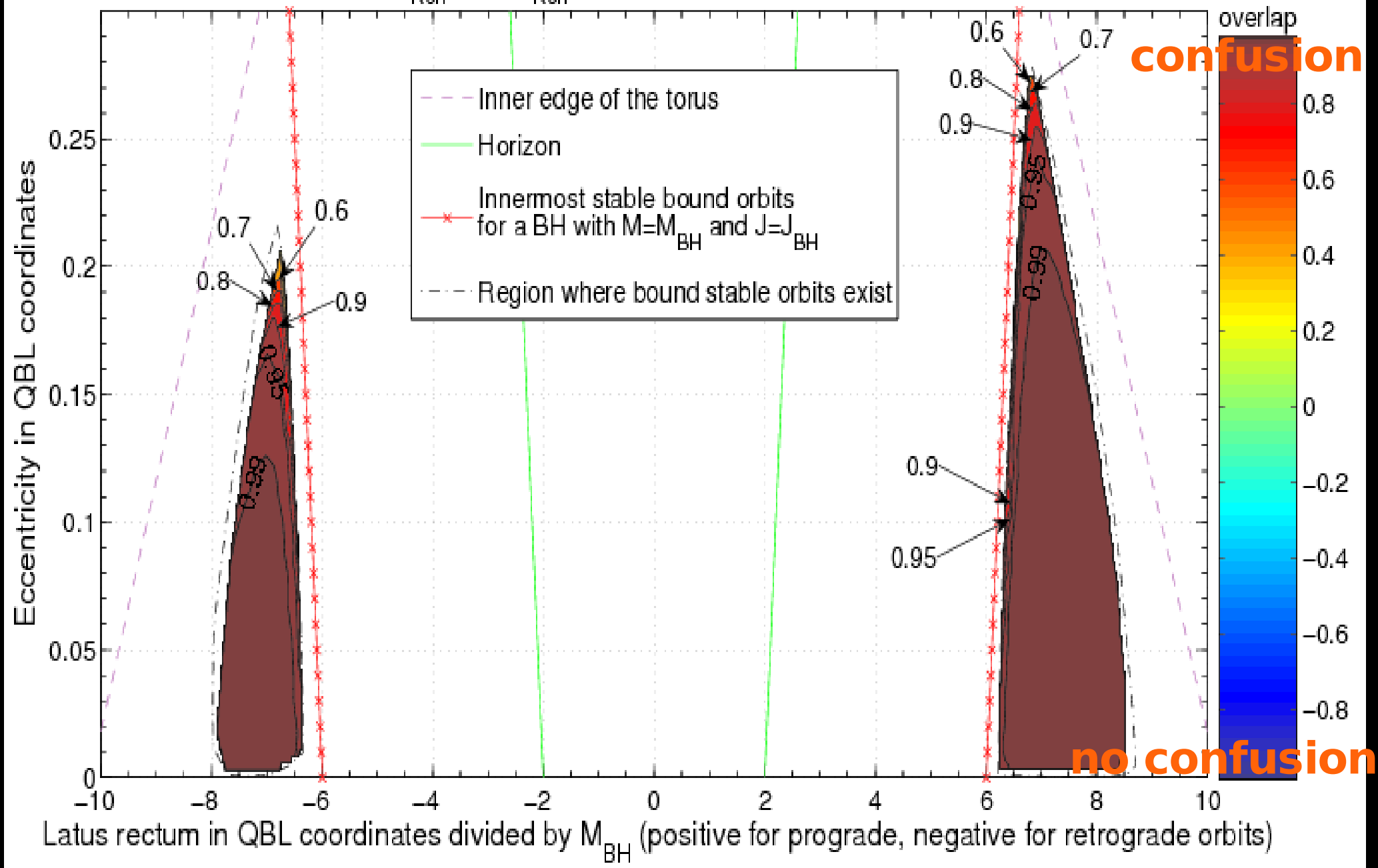


		Confusion in p and e	Confusion in a and M
Spacetime A	External Orbits	Only far from the system	Yes
	Internal Orbits	No	Yes
Spacetime B	External Orbits	Only far from the system	Yes

Same periods by changing latus rectum and eccentricity (J_{Kerr} and M_{Kerr} fixed)



Same periods by changing J_{Kerr} and M_{Kerr} (latus rectum and eccentricity in QBL coordinates fixed)



Conclusions and future work

- A confusion problem in the orbital parameters p and e arises only far from the BH+Torus system
- A confusion problem in the BH parameters a and M is present in most of the spacetime

Can this affect LISA's measurements?

Probably no. To confirm, we need

- More realistic tori (larger or lighter)
 - Inclusion of (approximate) radiation reaction
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