



# Phenomenological template bank for black hole coalescence waveforms

P. Ajith

*Albert Einstein Institute and University of Hannover*

*in collaboration with*

S. Babak, Y. Chen, B. Krishnan, B. Bruegmann, P. Diener,  
J. Gonzalez, M. Hannam, M. Hewitson, M. Koppitz, S. Husa, D. Pollney,  
L. Rezzolla, A. M. Sintes, U. Sperhake and J. Thornburg

16 APRIL 2007 **UNIVERSITY OF JENA**



- Data analysis from coalescing compact binaries – state of the art.
- Motivation and plans – look for BH coalescences with a single template family.
- Constructing full coalescence waveforms – Matching PN and NR waveforms.
- A phenomenological waveform family.
- Laying down templates – the template bank.
- Sensitivity & astrophysical range of the search.



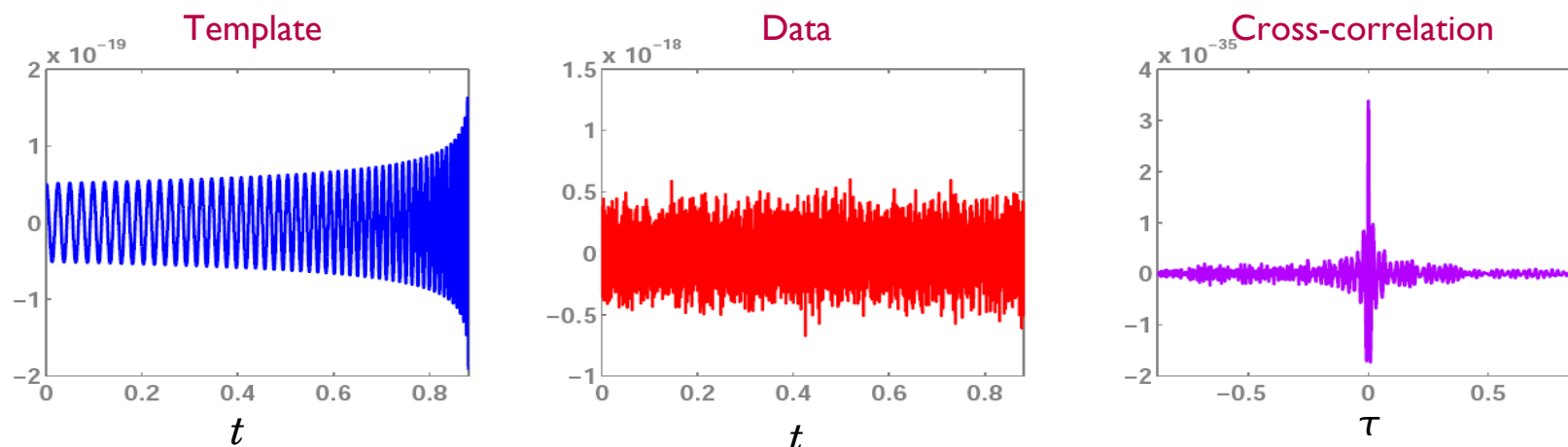
**This is a DA talk!**



# Data analysis for coalescing compact binaries

## ▪ **Inspiral phase**

- Physical (analytical) waveforms from PN theory. Also, detection template families (e.g. BCV) constructed to match with a number of theoretical waveforms.
- Since waveforms are *known*, the optimal filter to search these signals in the noise is *matched filter*, which essentially involves cross-correlating the data and the signal template.

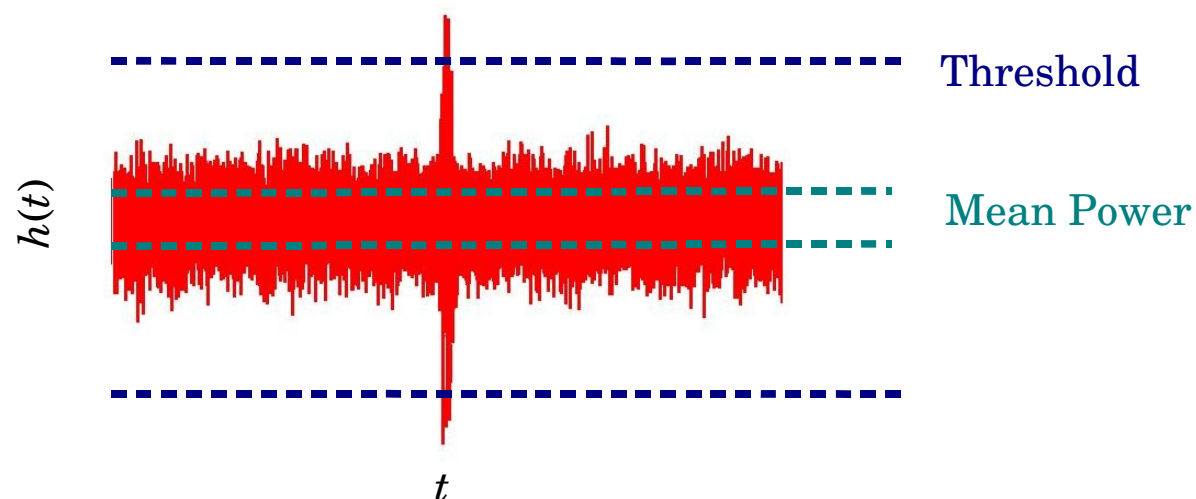




# Data analysis for coalescing compact binaries

- **Merger phase**

- Waveforms were not known until recently.
- Searches use *excess-power* techniques, which are tuned to detect large morphologies of waveforms. Look for short-lived excitations of power in the data, which are less-likely to be associated with the background noise distribution – a non-optimal filter.

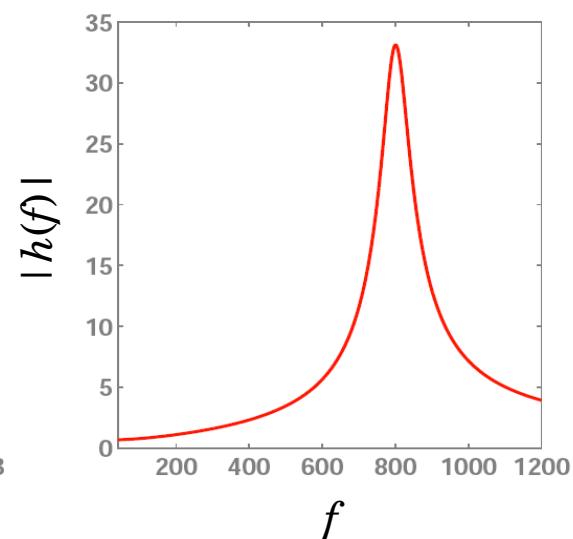
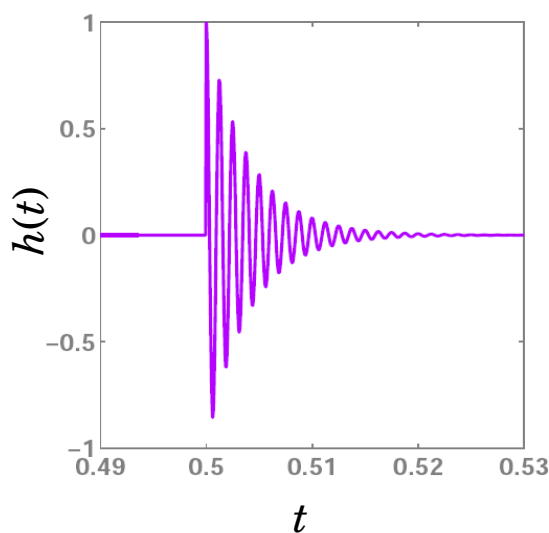




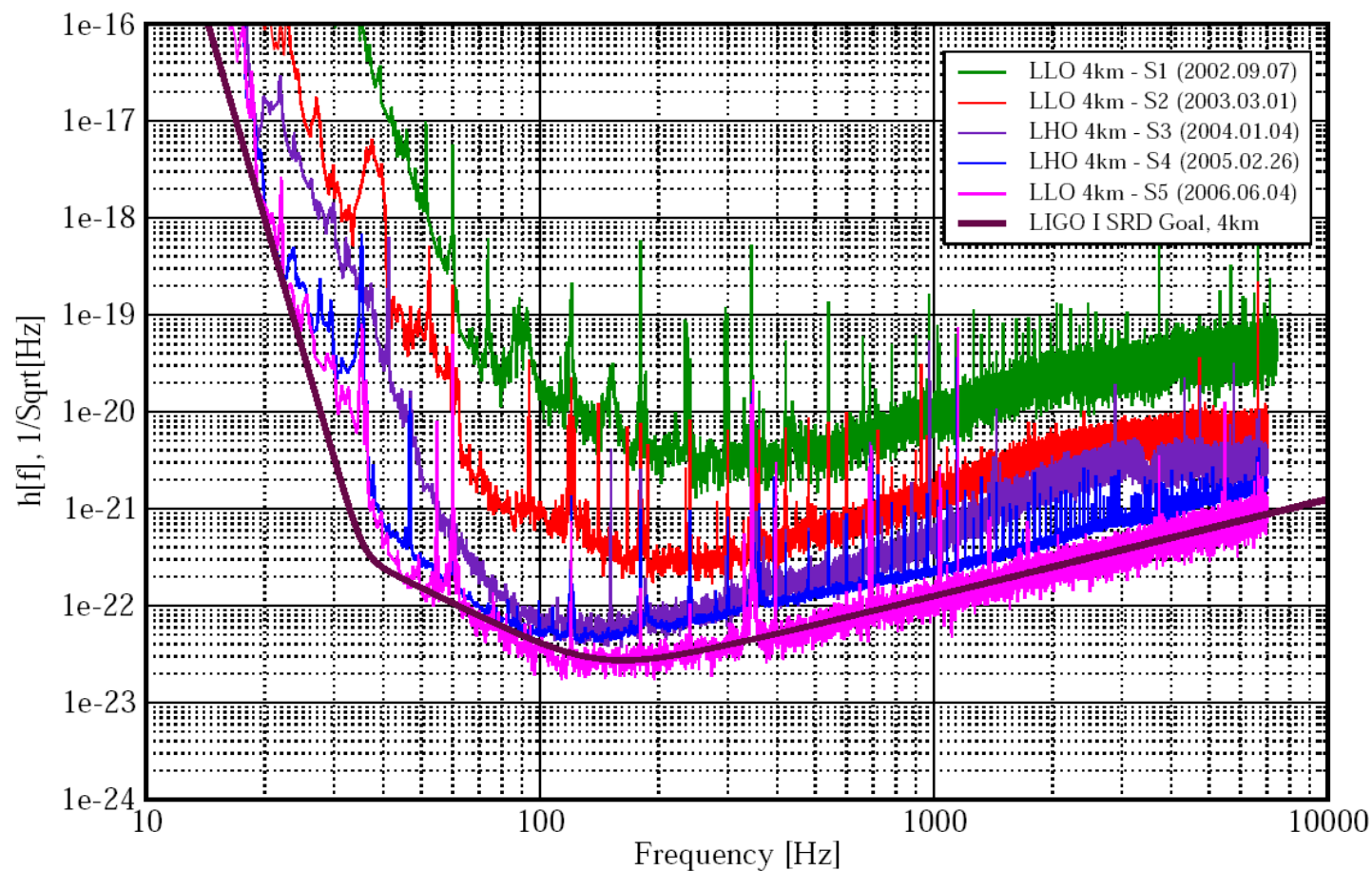
# Data analysis for coalescing compact binaries

- **Ring-down phase**

- Waveforms are exponentially-damped sinusoids, parametrized by the mass and spin of the BH.
- Since waveforms are known, matched filtering can be used to search for the signals.



# Motivation of the work



The ongoing science run of large-scale interferometers  
at the design sensitivity



## Motivation of the work

---

- Recent progress in Numerical Relativity in solving the binary BH problem.
- Gravitational waveforms from the non-perturbative merger phase can also be computed.
- Finally allows us to coherently search for all three stages (inspiral, merger and ring-down) of the binary BH coalescence.



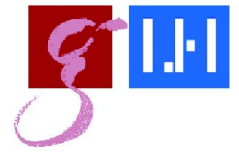
## A single template bank for BH coalescence

- Coherently search for all three stages of the BH coalescence signals (non-spinning binaries) using a single template bank.
- But the high computational cost makes it infeasible to generate enough numerical waveforms to densely cover the entire parameter space to be searched over using matched filtering technique.



# A single template bank for BH coalescence

- **Issue** How to construct a bank of templates?
  - Too expensive to compute a bank of NR waveforms dense enough in the  $(M, \eta)$  parameter space.
  - A phenomenological template bank (with parametrized waveforms) which has very good overlap with the 'target signals'.
- **Issue** How to construct the 'target' waveforms?
  - Need waveforms containing all three stages of the BH coalescence - too expensive to (numerically) evolve the binary from very large separations.
  - Match PN inspiral waveforms with NR (merger+ring-down) waveforms in a region where both calculations are valid.



## Matching PN and NR waveforms

- Minimize the 'distance' between PN and NR waveforms over a matching region (a few cycles long), thus construct hybrid waveforms.

$$\delta = \min_{\boldsymbol{\mu}, a} \left[ \sum_{i=+, \times} \int_{t_1}^{t_2} [h_i^{\text{PN}}(t, \boldsymbol{\mu}) - a h_i^{\text{NR}}(t, \boldsymbol{\nu})]^2 dt \right]$$

- Minimisation is carried out over the parameters  $\boldsymbol{\mu} = \{t_0, \phi_0, M, \eta\}$  and an amplitude scaling factor  $a$ .

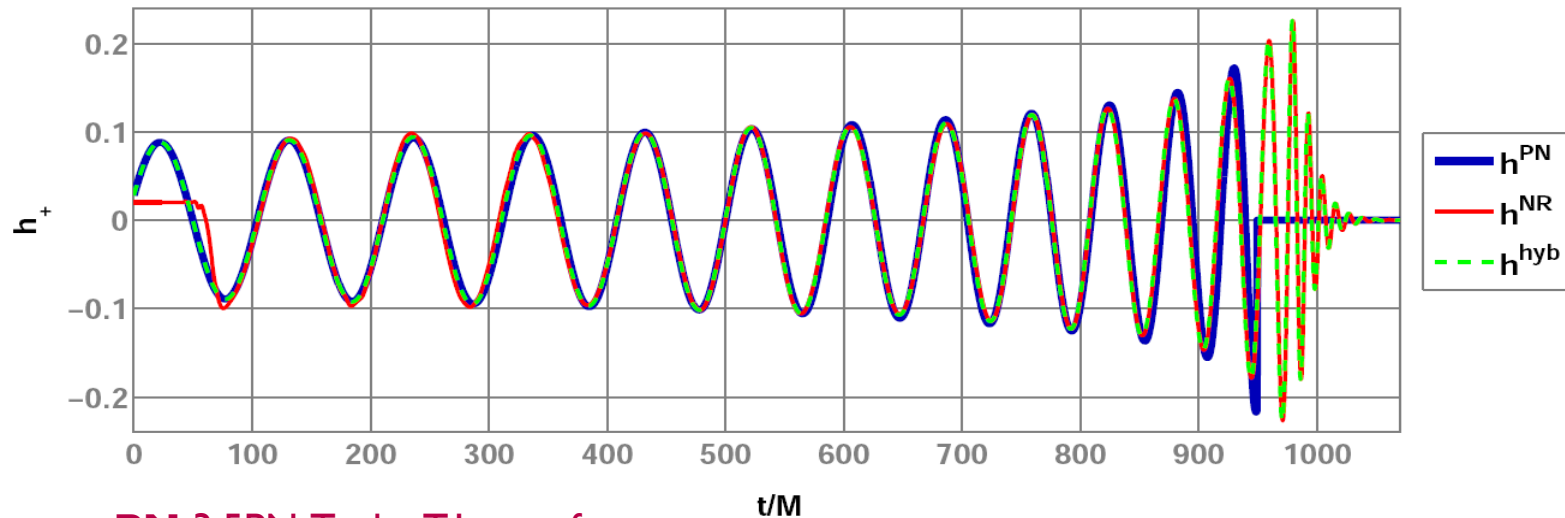
## Matching PN and NR waveforms

- Minimize the 'distance' between PN and NR waveforms over a matching region (a few cycles long), thus construct hybrid waveforms.

$$\delta = \min_{\mu, a} \left[ \sum_{i=+, \times} \int_{t_1}^{t_2} [h_i^{\text{PN}}(t, \mu) - a h_i^{\text{NR}}(t, \nu)]^2 dt \right]$$

- Minimisation is carried out over the parameters  $\mu = \{t_0, \phi_0, M, \eta\}$  and an amplitude scaling factor  $a$ .

# Constructing hybrid waveforms

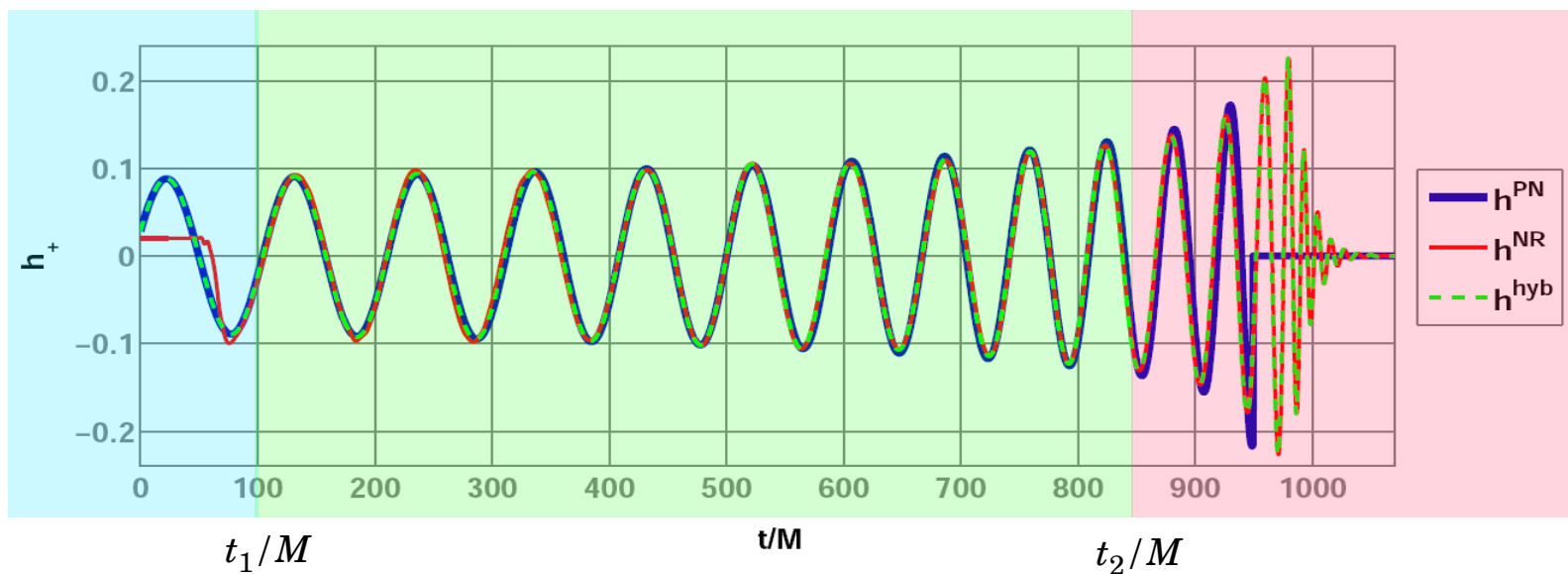


**PN** 3.5PN TaylorT1 waveform  
**NR** AEI equal-mass simulation

- Combine NR waveforms with the best-matched PN waveforms

$$h_{+, \times}^{\text{hyb}}(t, \boldsymbol{\nu}) = \begin{cases} h_{+, \times}^{\text{PN}}(t, \boldsymbol{\mu}_0) & \text{if } t < t_1 \\ a_0 \tau h_{+, \times}^{\text{NR}}(t, \boldsymbol{\nu}) + (1 - \tau) h_{+, \times}^{\text{PN}}(t, \boldsymbol{\mu}_0) & \text{if } t_1 \leq t < t_2 \\ a_0 h_{+, \times}^{\text{NR}}(t, \boldsymbol{\nu}) & \text{if } t_2 \leq t \end{cases}$$

# Constructing hybrid waveforms

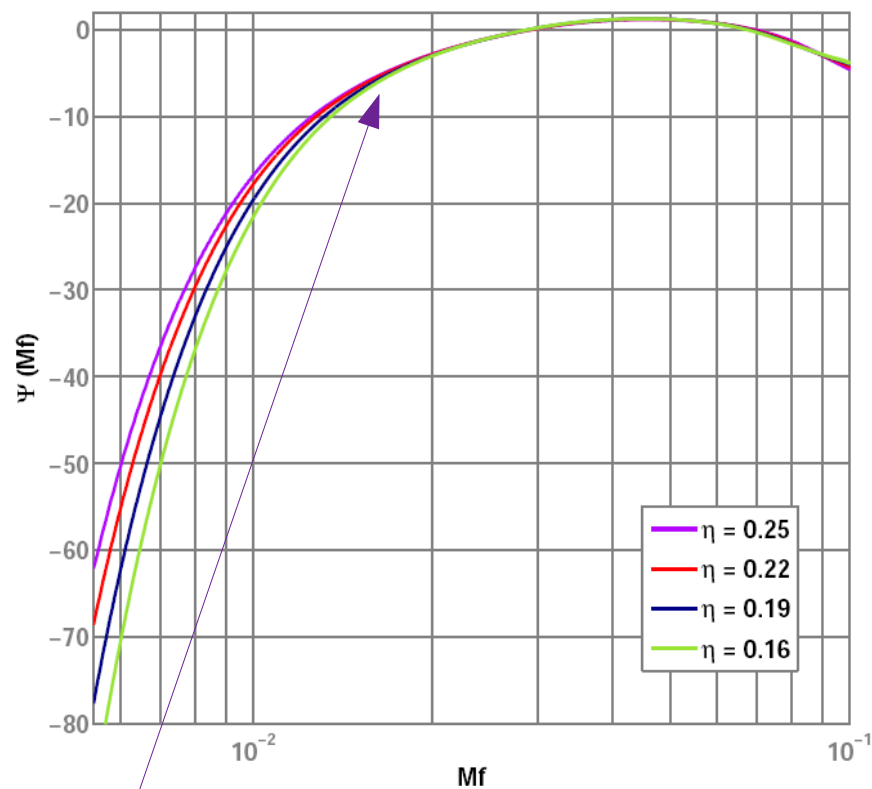
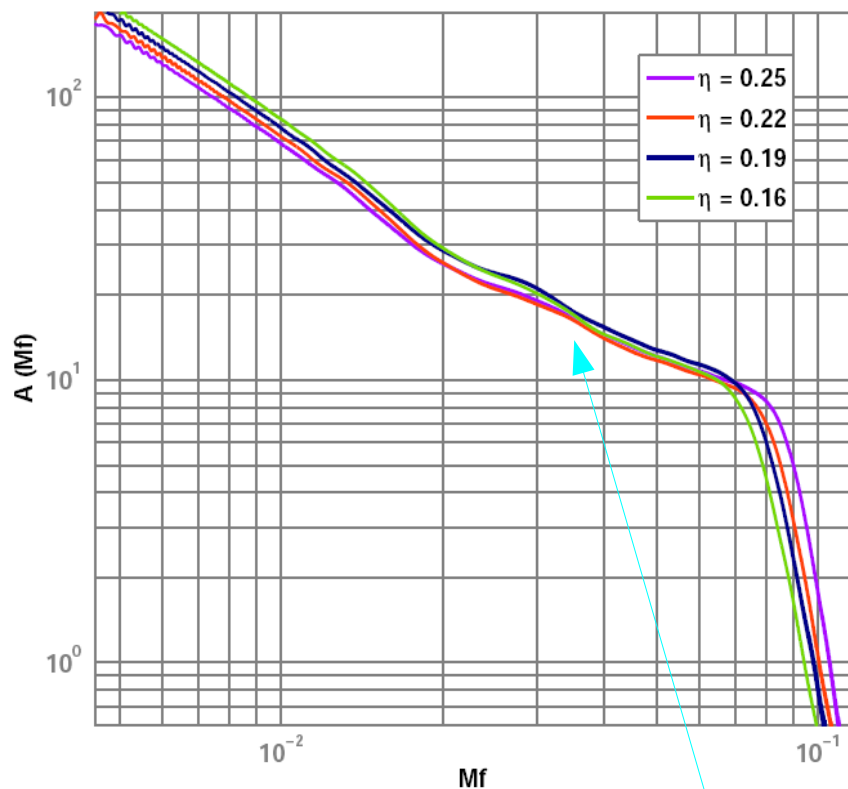


'Best-matched' parameters

$$h_{+, \times}^{\text{hyb}}(t, \nu) = \begin{cases} h_{+, \times}^{\text{PN}}(t, \mu_0) & \text{if } t < t_1 \\ a_0 \tau h_{+, \times}^{\text{NR}}(t, \nu) + (1 - \tau) h_{+, \times}^{\text{PN}}(t, \mu_0) & \text{if } t_1 \leq t < t_2 \\ a_0 h_{+, \times}^{\text{NR}}(t, \nu) & \text{if } t_2 \leq t \end{cases}$$

Linearly increasing weighting function such that  $\tau(t_1) = 0$  &  $\tau(t_2) = 1$

# Hybrid waveforms in the Fourier domain

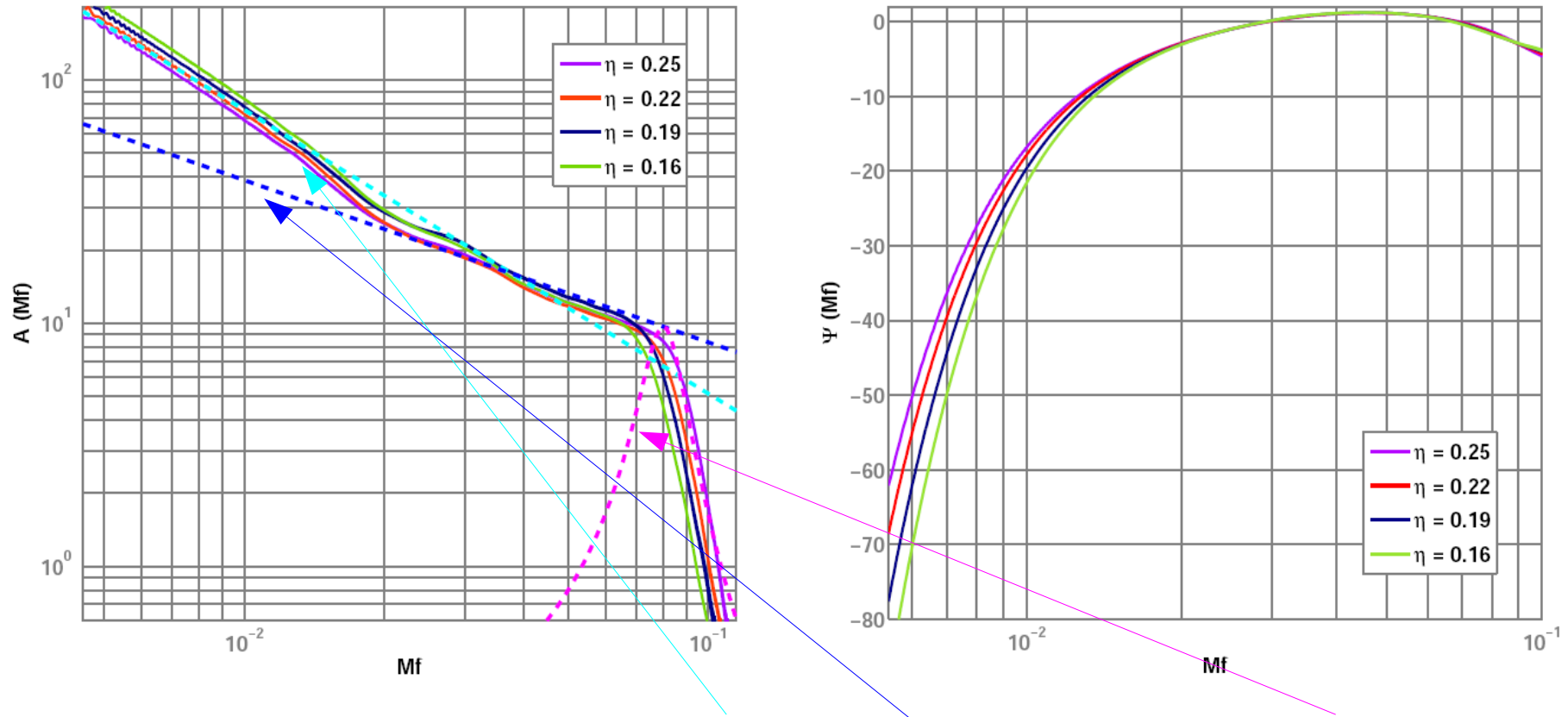


- Fourier domain magnitude and phase of hybrid waveforms.

**PN** 3.5PN TaylorT1 waveform

**NR** Jena unequal-mass simulations

# A phenomenological parametrization



- **Magnitude** by two power-laws ( $f^{-7/6}$  and  $f^{-2/3}$ ) and a Lorentzian  $L(f_{\text{ring}}, \sigma)$ .
- **Phase** Expansion in powers of  $f$  (motivation from the SPA).

## Phenomenological waveforms

- The parametrised phenomenological waveform is written as

$$u(f) = \mathcal{A}_{\text{eff}}(f) e^{i\Psi_{\text{eff}}(f)}$$

where

$$\mathcal{A}_{\text{eff}}(f) = \begin{cases} (f/f_{\text{merg}})^{-7/6} & \text{if } f < f_{\text{merg}} \\ (f/f_{\text{merg}})^{-2/3} & \text{if } f_{\text{merg}} \leq f < f_{\text{ring}} \\ w \mathcal{L}(f, f_{\text{ring}}, \sigma) & \text{if } f_{\text{ring}} \leq f < f_{\text{cut}} \end{cases}$$

$$\begin{aligned} \Psi_{\text{eff}}(f) = & 2\pi f t_0 + \phi_0 + \psi_0 f^{-5/3} + \psi_2 f^{-1} \\ & + \psi_3 f^{-2/3} + \psi_4 f^{-1/3} + \psi_6 f^{1/3}. \end{aligned}$$



# Phenomenological waveforms

- The parametrised phenomenological waveform is written as

$$u(f) = \mathcal{A}_{\text{eff}}(f) e^{i\Psi_{\text{eff}}(f)}$$

where

$$\mathcal{A}_{\text{eff}}(f) = \begin{cases} (f/f_{\text{merg}})^{-7/6} & \text{if } f < f_{\text{merg}} \\ (f/f_{\text{merg}})^{-2/3} & \text{if } f_{\text{merg}} \leq f < f_{\text{ring}} \\ w \mathcal{L}(f, f_{\text{ring}}, \sigma) & \text{if } f_{\text{ring}} \leq f < f_{\text{cut}} \end{cases}$$

Transition frequencies

'Spread' of the Lorentzian

Cutoff frequency

$$\Psi_{\text{eff}}(f) = 2\pi f t_0 + \phi_0 + \psi_0 f^{-5/3} + \psi_2 f^{-1} + \psi_3 f^{-2/3} + \psi_4 f^{-1/3} + \psi_6 f^{1/3}.$$

Phasing coefficients



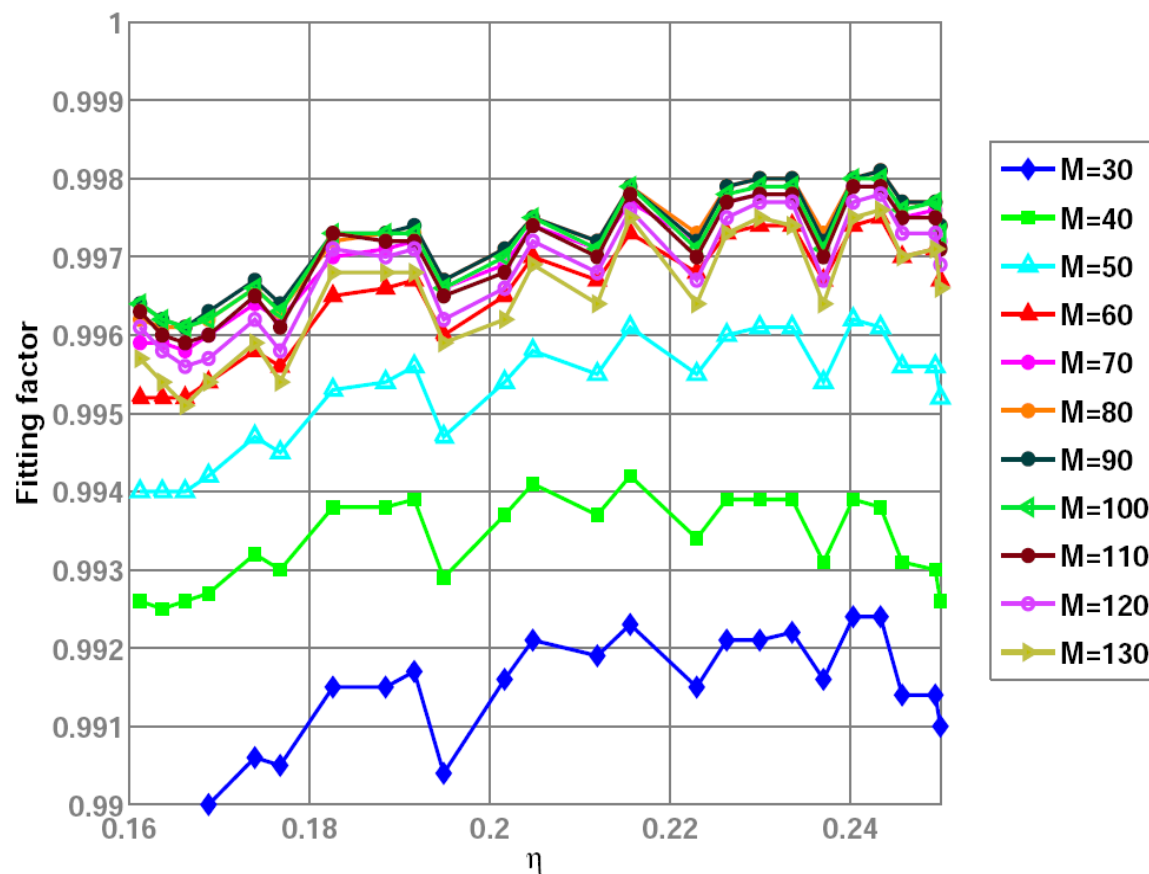
## Characterising the quality of the phenomenological waveforms

- **Fitting factor** Ratio of the SNR that can be achieved with suboptimal template to the SNR obtained with an optimal template -- overlap between the (normalized) template and signal, maximized over the parameters of the template.

$$\text{FF}(h(f, \boldsymbol{\mu}), u(f, \boldsymbol{\alpha})) = \frac{\max_{\boldsymbol{\alpha}} \langle h(f, \boldsymbol{\mu}), u(f, \boldsymbol{\alpha}) \rangle}{\sqrt{\langle h(f, \boldsymbol{\mu}), h(f, \boldsymbol{\mu}) \rangle} \sqrt{\langle u(f, \boldsymbol{\alpha}), u(f, \boldsymbol{\alpha}) \rangle}},$$

# Fitting factors with the hybrid waveforms

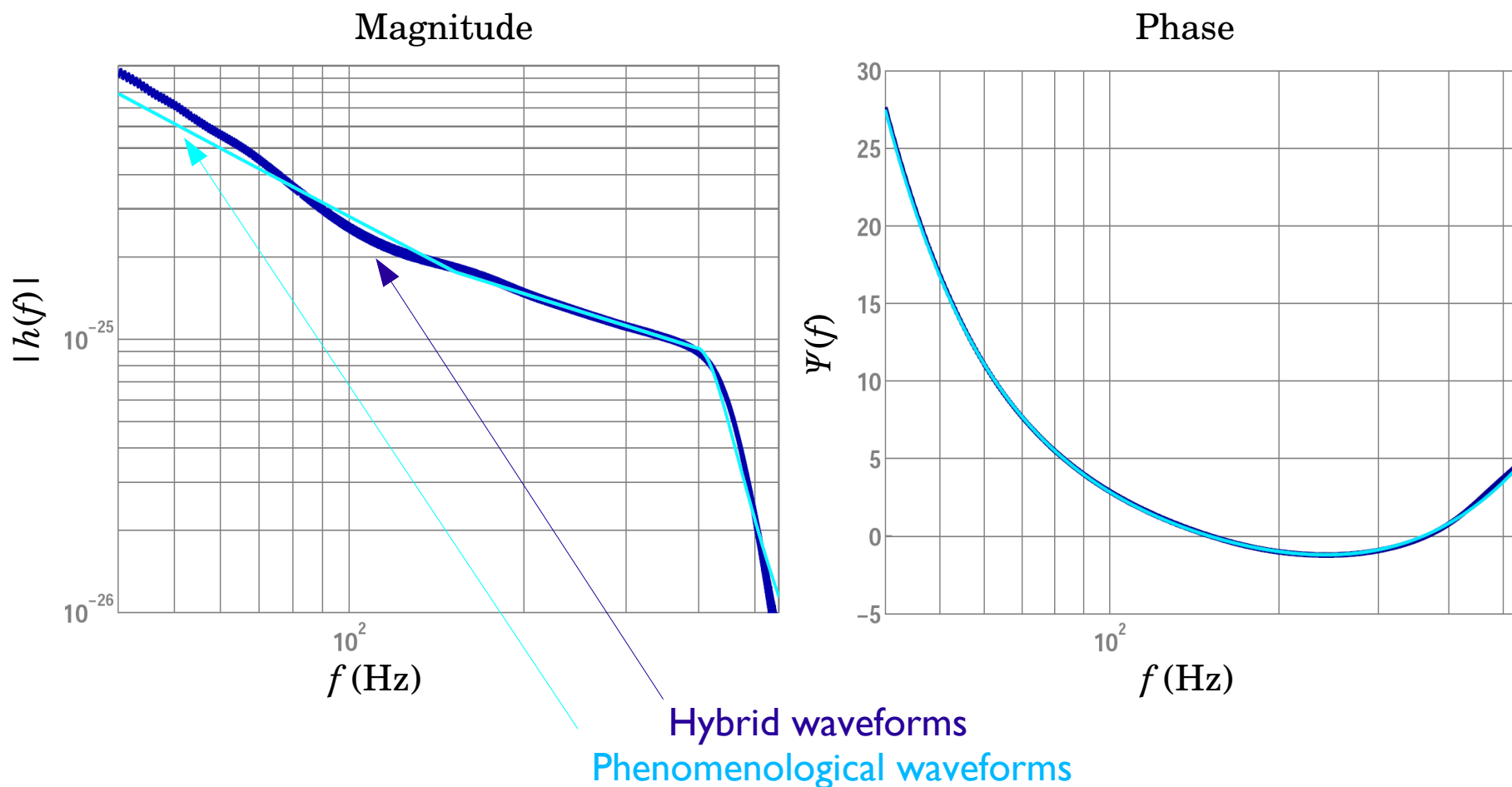
- Hybrid waveforms**  
 Jena UM simulations  
 $(0.16 \leq \eta \leq 0.25)$  +  
 3.5PN inspiral.
- Noise spectrum**  
 Initial LIGO design.





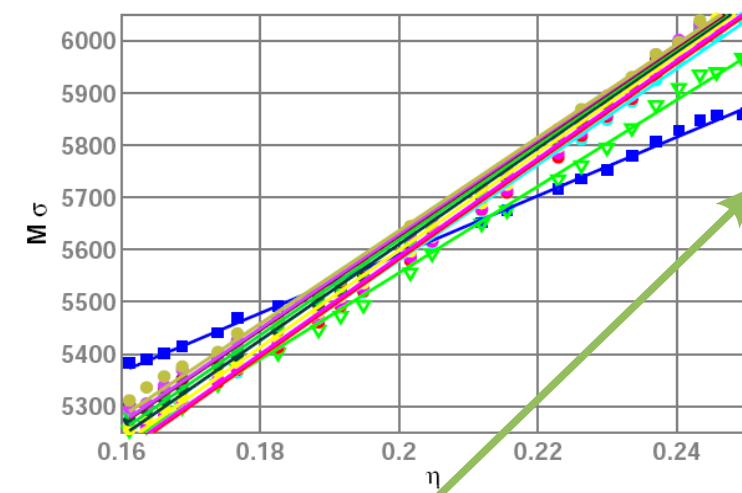
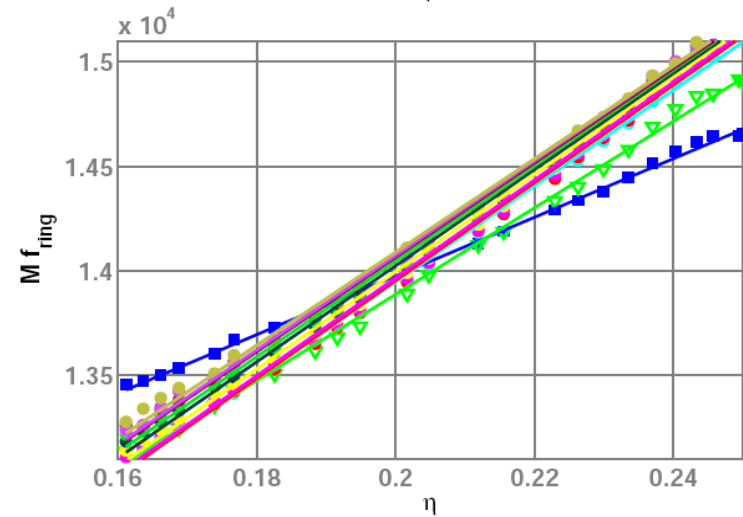
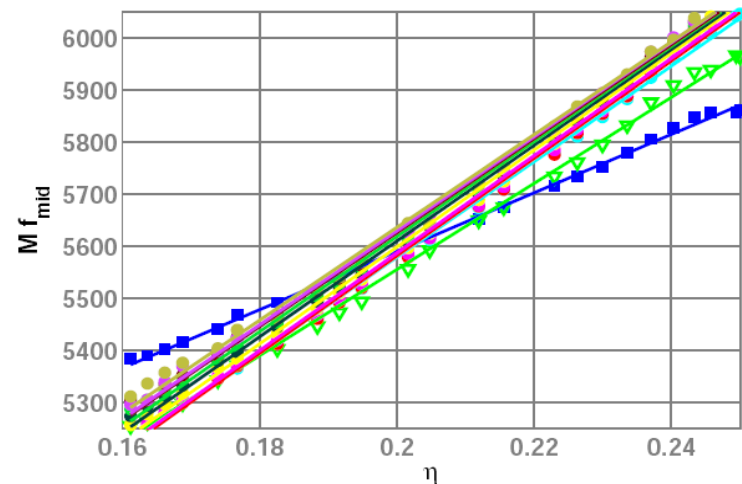
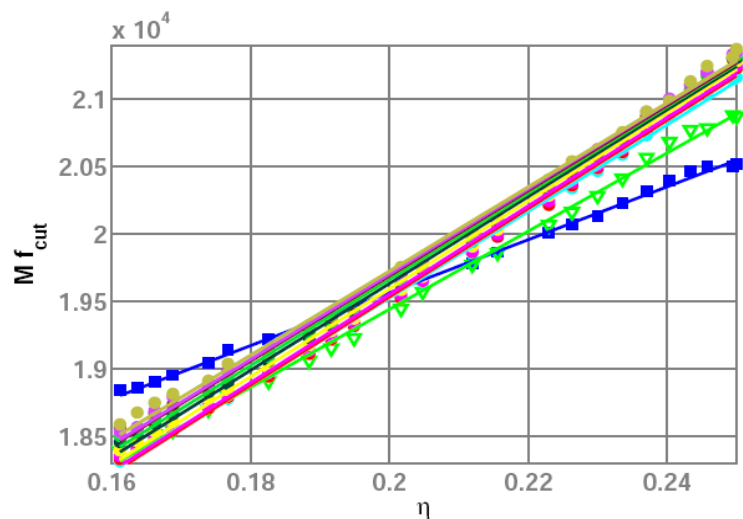
# Hybrid waveforms and the 'best-matched' templates

- $M = 40 M_{\odot}$ ,  $\eta = 0.25$  system, Initial LIGO noise spectrum.

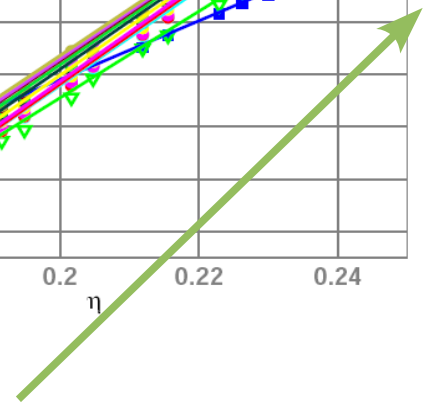


# From phenomenological parameters to physical parameters

Best-matched phenomenological parameters



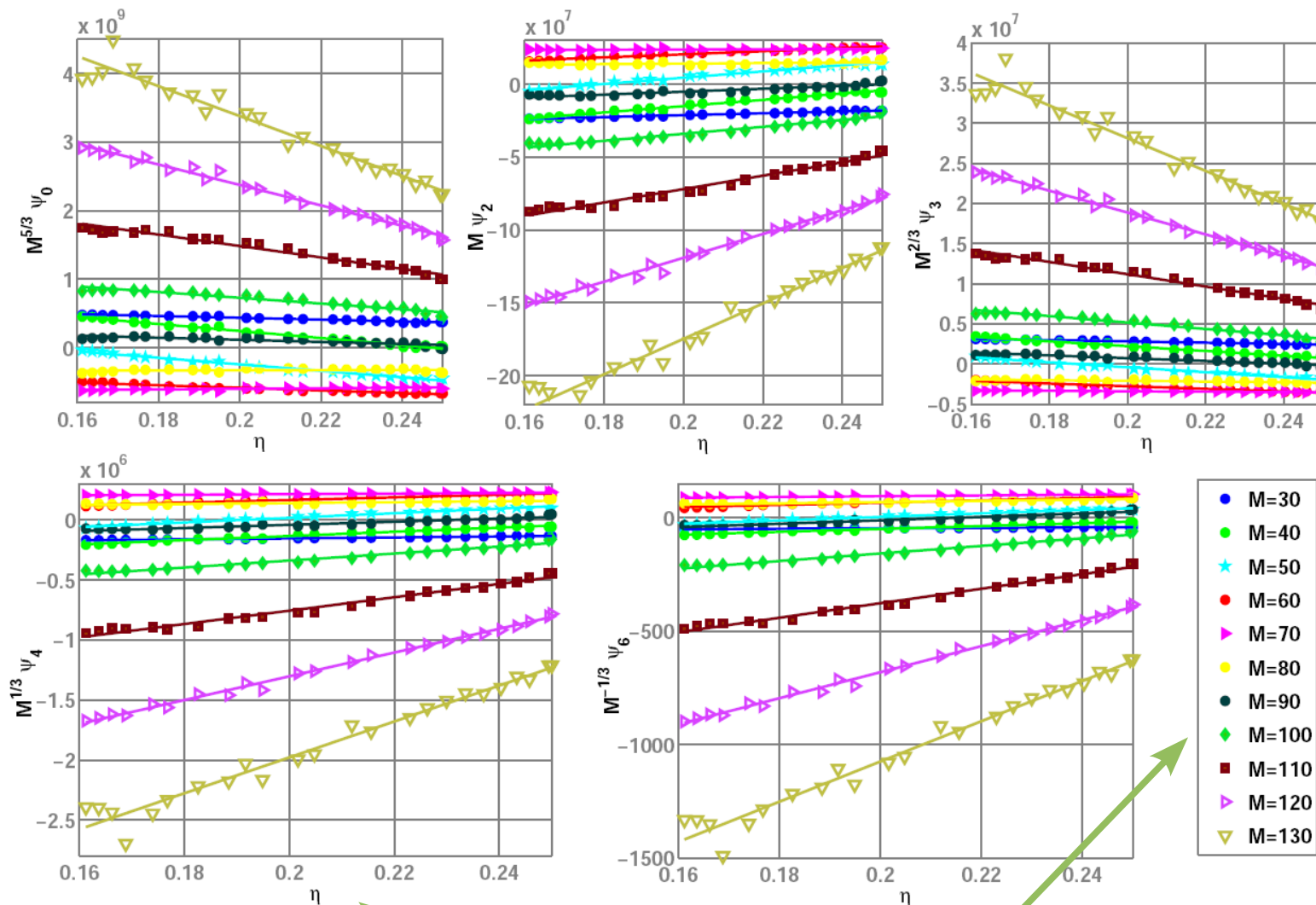
- M=30
- ▼ M=40
- M=50
- M=60
- M=70
- M=80
- M=90
- M=100
- M=110
- M=120
- M=130



Physical parameters of the binary

# From phenomenological parameters to physical parameters

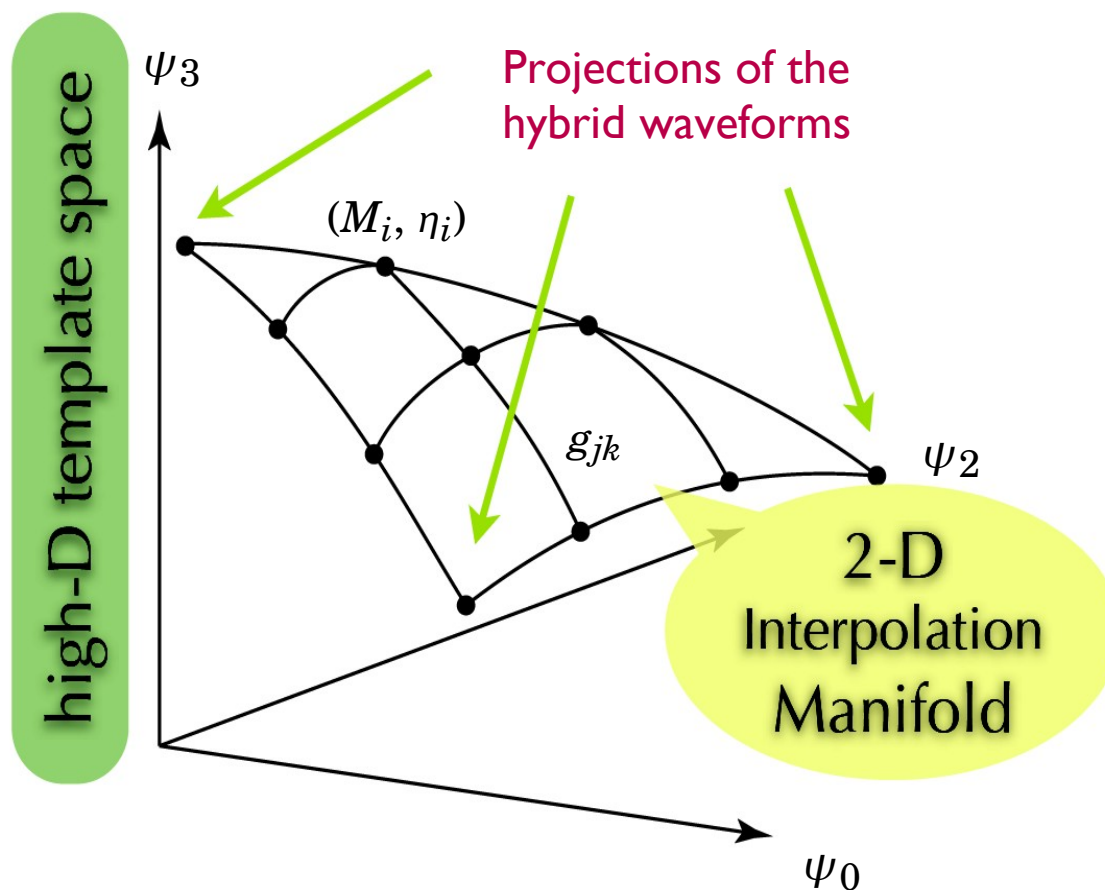
Best-matched phenomenological parameters



Physical parameters of the binary

# Template bank

- Re-parametrize the templates in terms of  $M$  and  $\eta$ . The template family is a two dimensional manifold embedded in a higher dimensional space.
- A metric can be defined on this manifold, which gives the notion of distance. This can be used to lay down templates allowing a given mismatch between neighboring templates.

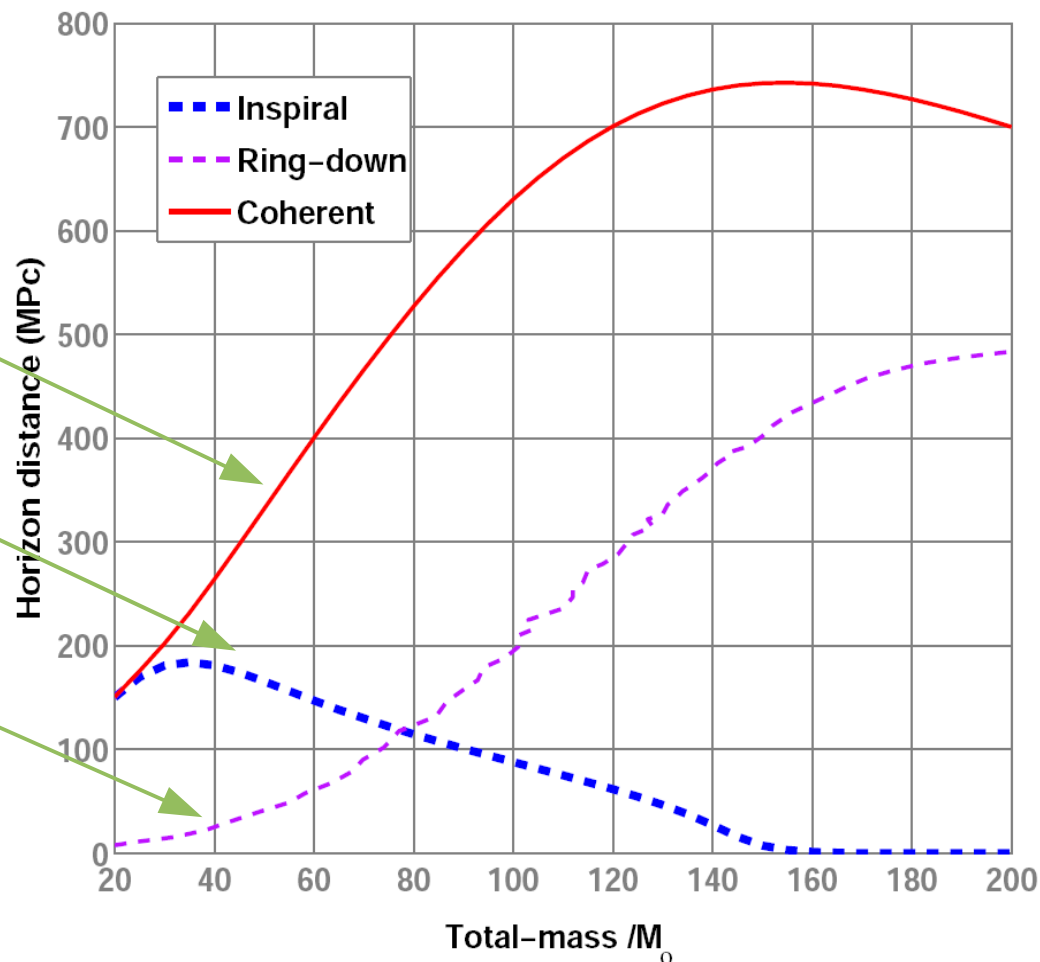


# Sensitivity of the search

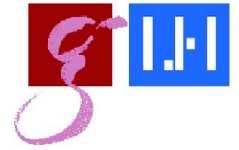
Using inspiral+merger+ring-down templates

Using standard PN templates (truncated at ISCO)

Using ring-down templates



- Effective distance to optimally-oriented systems which can produce an optimal SNR of 8 at Initial LIGO.



## Summary

- Recent progress in Numerical Relativity in modelling the non-perturbative merger phase of the binary BH coalescence problem.
- Constructed a set of hybrid waveforms by matching PN and NR predictions – waveforms containing the inspiral, merger, and ring-down stages.
- Proposed a phenomenological waveform family which has very good overlaps with the hybrid waveforms. These waveforms can be parametrized in terms of the physical parameters of the system – a two-dimensional template bank (for non-spinning binaries).
- This template bank might enable us to extend the present inspiral searches to higher mass binary black hole systems – increased reach of the current generation of ground based detectors.